

Quantum fluctuations and the phase diagram of anisotropic dipolar magnets

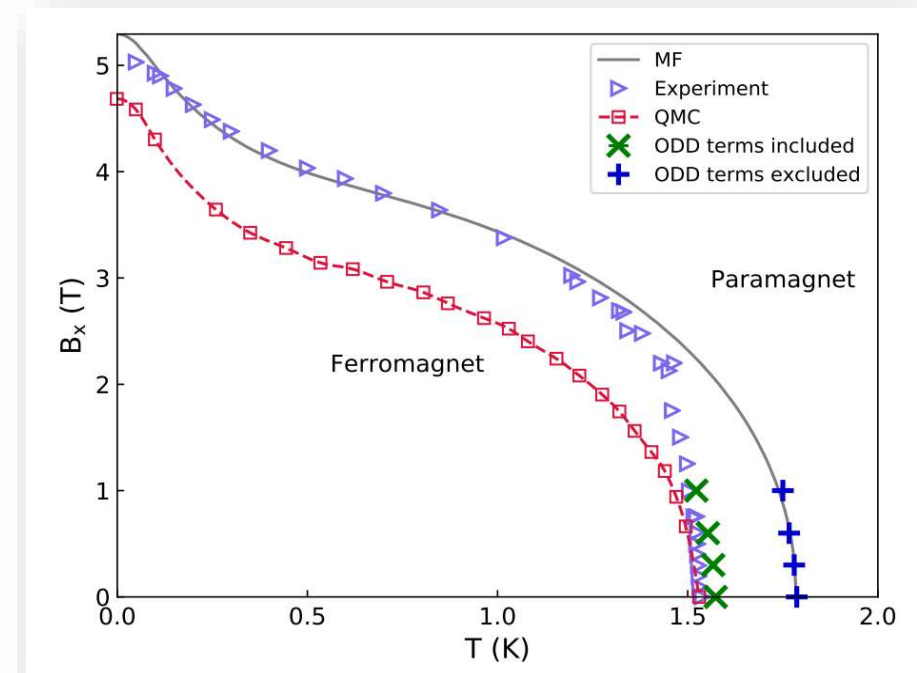
Correlated Days, Ein Gedi, 06/2022

Tomer Dollberg (BGU)

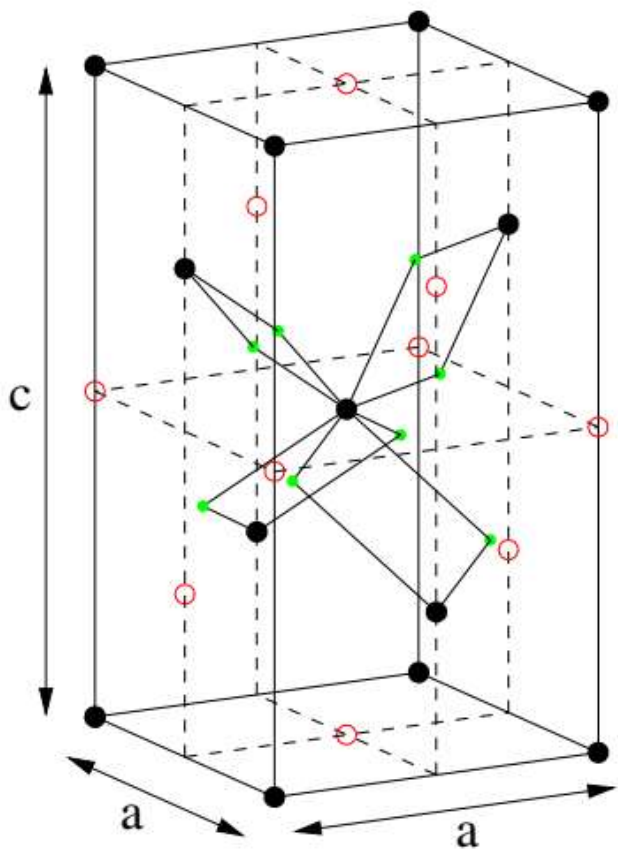
Moshe Schechter (BGU)

Juan Carlos Andresen (BGU)

TD, J. C. Andresen, and M. Schechter, Phys. Rev. B **105**, L180413 (2022)

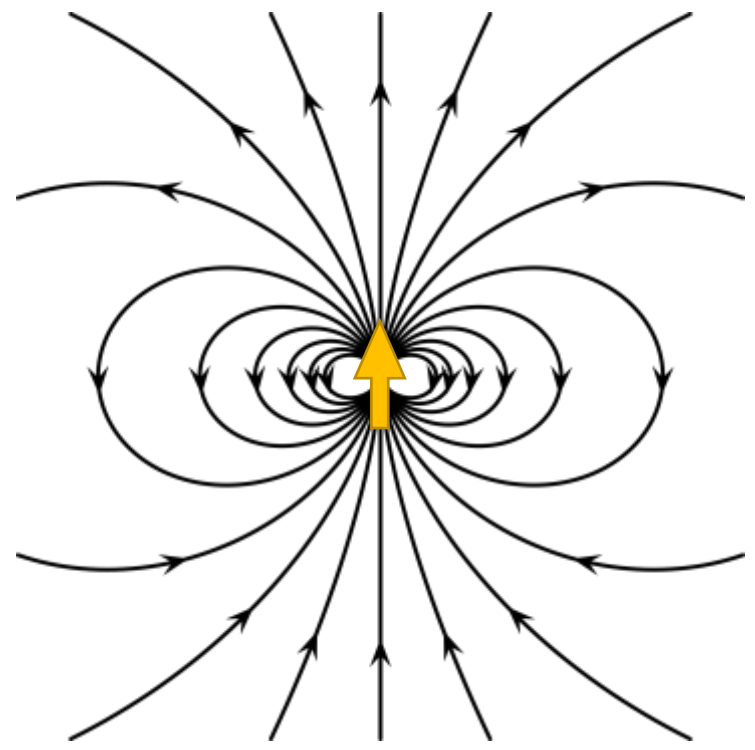


Introduction

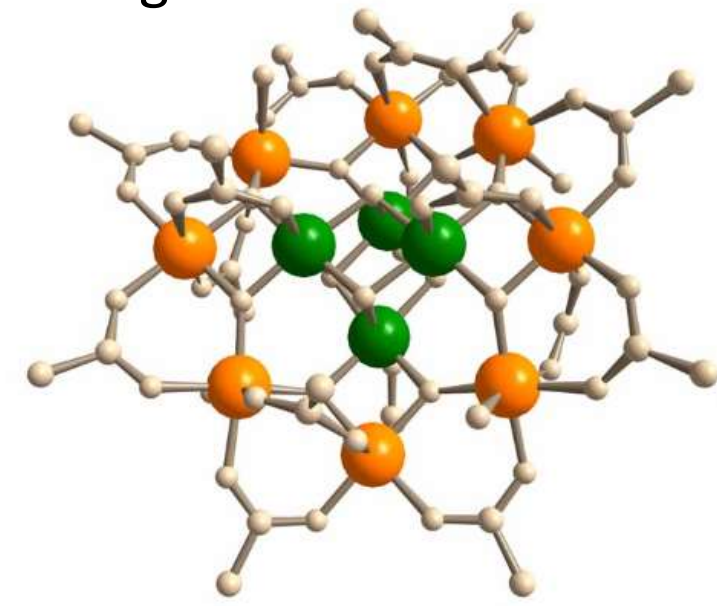


- Ho^{3+}
- Li^{+}
- F^{-}

Rare-earth magnetic insulators



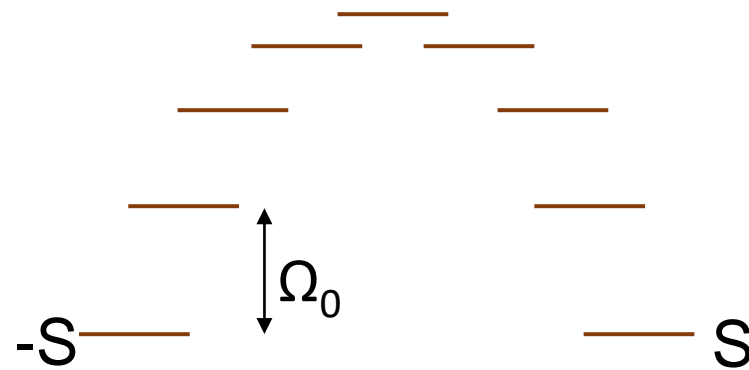
Single molecule magnets



Anisotropic dipolar systems

Magnetic insulators, large spin, strong lattice anisotropy, dominant dipolar interaction

$$H_{\text{cf}} = -D \sum_i S_i^z{}^2$$



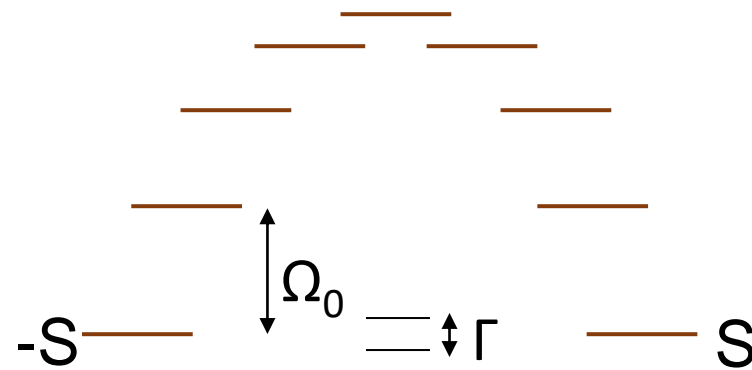
$$H_D = H_{\text{cf}} - \sum_{ij} V_{ij}^{\mu\nu} S_i^\mu S_j^\nu$$

$$H_{\text{Ising}} = - \sum_{i,j} J_{ij} \sigma_i^z \sigma_j^z$$

Anisotropic dipolar systems

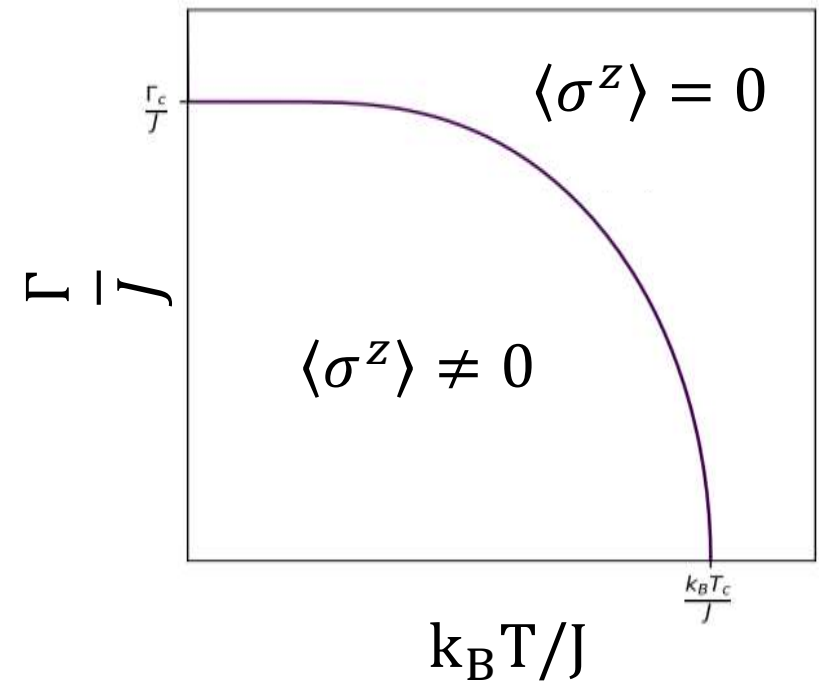
Magnetic insulators, large spin, strong lattice anisotropy, dominant dipolar interaction

$$H_{\text{cf}} = -D \sum_i S_i^z{}^2$$



$$H_D = H_{\text{cf}} - \sum_{ij} V_{ij}^{\mu\nu} S_i^\mu S_j^\nu - \Delta \sum_i S_i^x$$

$$H_{\text{Ising}} = - \sum_{i,j} J_{ij} \sigma_i^z \sigma_j^z - \Gamma \sum_i \sigma_i^x \quad \Gamma \propto \Delta^{2S}$$



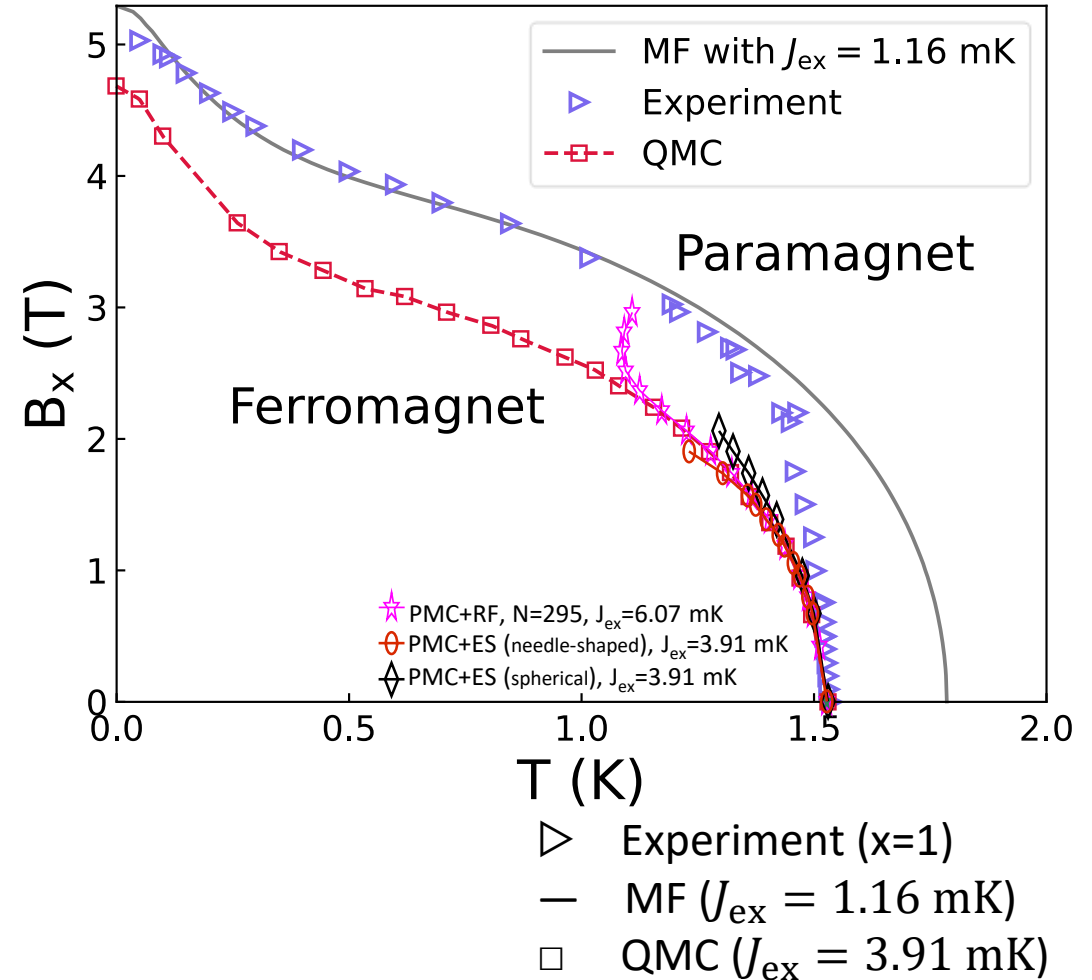
Mean field (MF) phase diagram of transverse field Ising model (TFIM)

LiHoF₄ – the archetypal Ising magnet?

- Full microscopic Hamiltonian ($J=8$, $I=7/2$):

$$H = \sum_i H_{\text{cf}}(\mathbf{J}_i) - g_L \mu_B \sum_i B_x J_i^x + \frac{1}{2} (g_L \mu_B)^2 \sum_{i \neq j} V_{ij}^{\mu\nu} J_i^\mu J_j^\nu + J_{\text{ex}} \sum_{\langle i,j \rangle} \mathbf{J}_i \cdot \mathbf{J}_j + A \sum_i (\mathbf{I}_i \cdot \mathbf{J}_i)$$

- The only free parameter, J_{ex} , can be used to tune $T_c(B_x = 0)$
- Something missing!



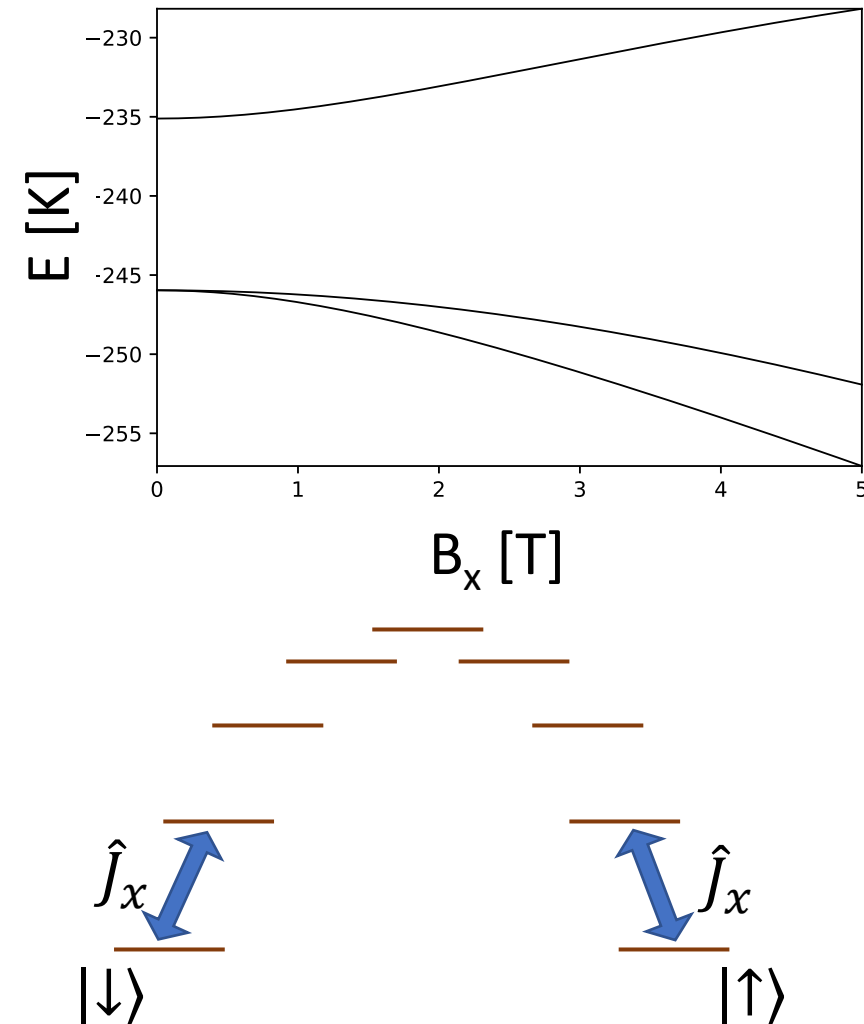
Chakraborty et al., PRB 70, 144411 (2004)

Tabei et al., PRB 78, 184408 (2008)

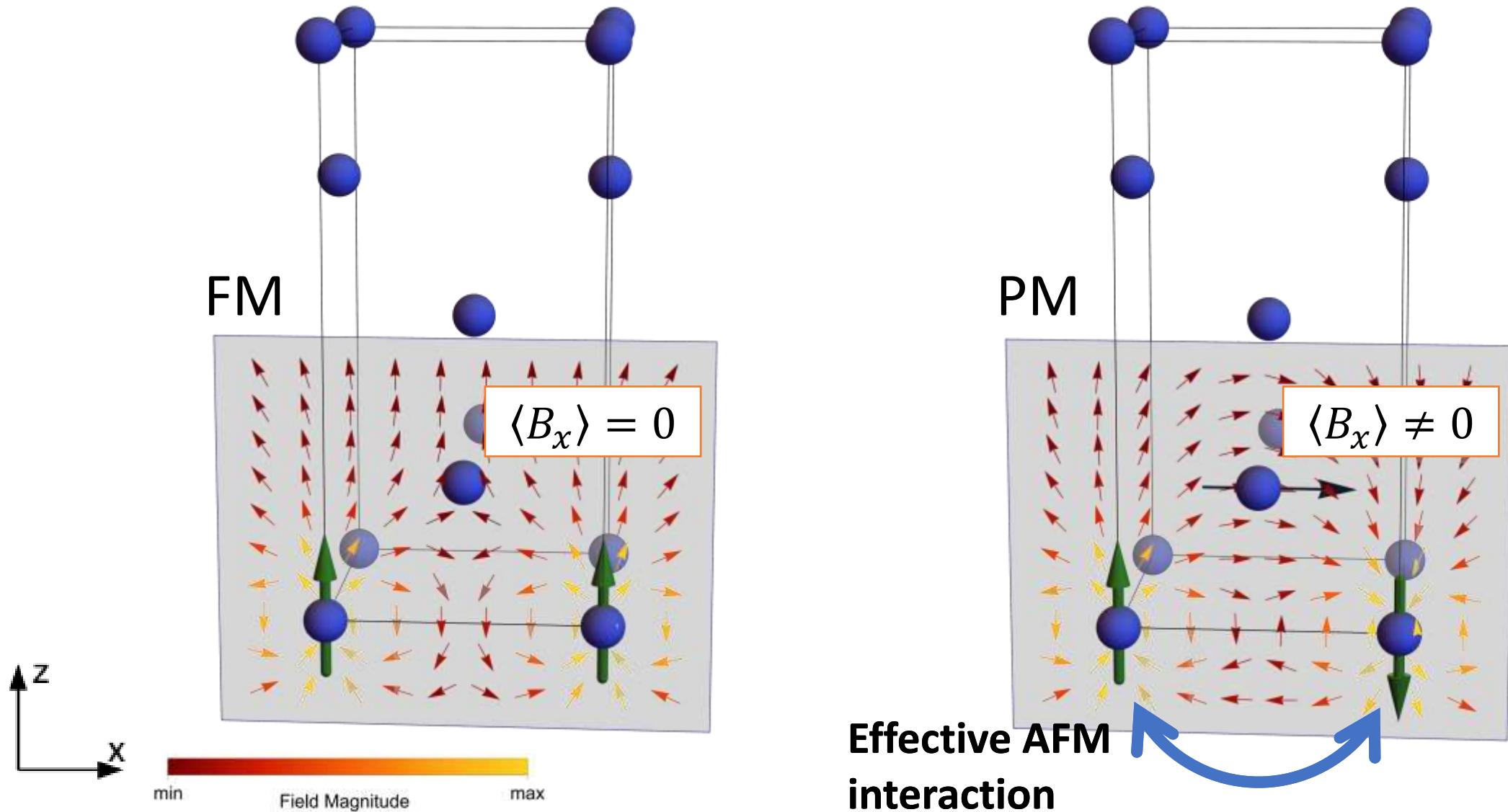
Babkevich et al., PRB 94, 174443 (2016)

Off-diagonal dipolar (ODD) terms

- Terms of the form $V_{ij}^{zx} \langle J_i^z \rangle J_j^x$
 - Effectively act as internal transverse magnetic fields
 - Decrease the energy of spin j regardless of its state
 - Neglected in TFIM due to projection
 - Vanish due to symmetry *in FM phase*
- Energetically disfavor ferromagnetic order



ODD terms and the phase diagram - example



Studied Hamiltonian – Monte Carlo Simulation

- Starting from the full microscopic Hamiltonian we keep ODD terms
- Approximation: for each spin, each of the other N-1 spins are projected not *one* of their single-site states — neglecting quantum many body effects
- Effective Hamiltonian

- Hybrid quantum-classical Monte Carlo

$$H_{\text{eff}} = \sum_i H_{\text{cf}}(\mathbf{J}_i) - g_L \mu_B \sum_i \langle \mathbf{B}_i \rangle \cdot \mathbf{J}_i$$

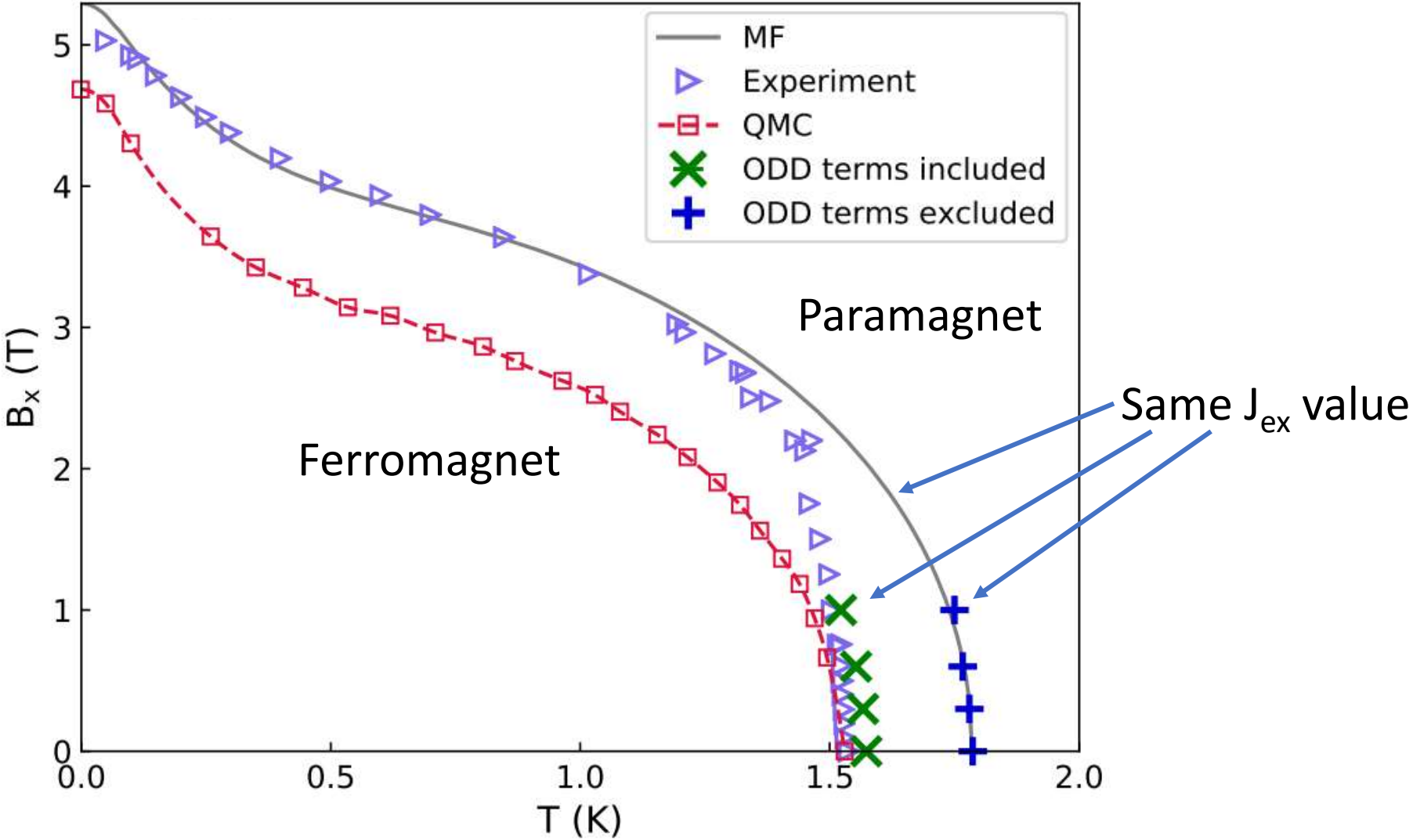
$$B_i^x = B_x - g_L \mu_B \sum_{j \neq i} V_{ij}^{zx} J_j^z$$

$$B_i^y = -g_L \mu_B \sum_{j \neq i} V_{ij}^{zy} J_j^z$$

$$B_i^z = -\frac{1}{2} g_L \mu_B \sum_{j \neq i} V_{ij}^{zz} J_j^z - \frac{J_{\text{ex}}}{2g_L \mu_B} \sum_{j \in \text{N.N.}} J_j^z$$

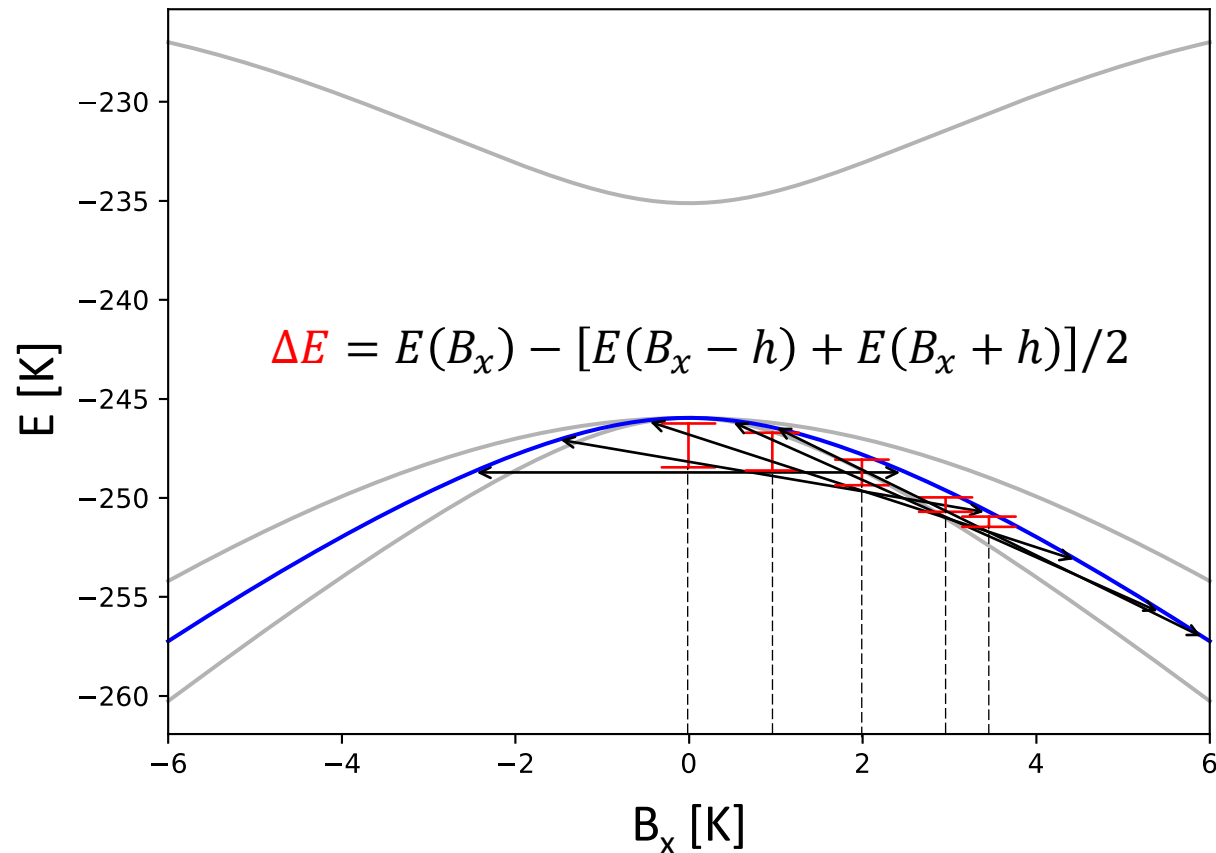
$$(J_{\text{ex}} = 1.16 \text{ mK})$$

Numerical results

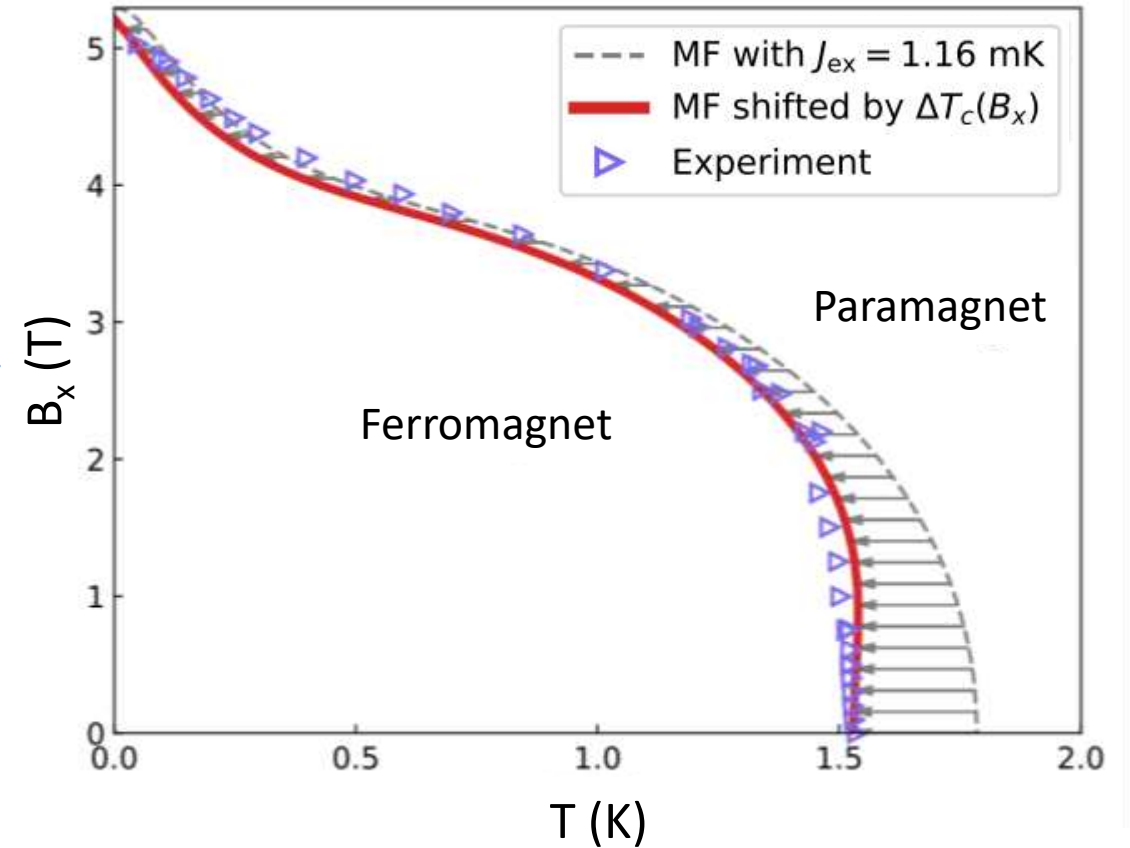


Can it explain the full phase boundary line?

ΔE is an estimation of energy difference between the PM and FM phases



Further, assume $\Delta T_c \propto \Delta E$



Conclusions

- Quantum fluctuations induced by off-diagonal dipolar terms affect even the classical phase transition of dipolar Ising systems
- The effect diminishes with increasing external B_x
- Description of anisotropic dipolar systems by the Ising model essentially insufficient

Thank you!