

E5651: Ising in interaction with lattice gas

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The problem:

Consider a one dimensional Ising model of spins $\sigma_i = \pm 1, i = 1, 2, 3, \dots, N$ and circular boundary condition ($\sigma_{N+1} = \sigma_1$). Between each two spins there is a site for an additional atom, which if present decreases the coupling J to $J(1 - \lambda)$. The Hamiltonian is then

$$H = -J \sum_{i=1}^N \sigma_i \sigma_{i+1} (1 - \lambda n_i)$$

where $n_i = 0$ or 1 and there are $N' = \sum_{i=1}^N n_i$ atoms ($N' < N$), i.e. on average $\langle n_i \rangle = N'/N$.

(1) Evaluate the partition sum by allowing all configurations of spins and of atoms. Calculate the free energy.

(2) If the atoms are stationary impurities one needs to evaluate the free energy F for some random configuration of the atoms and then average F over all configurations. Evaluate the average F . Find the entropy difference between the result calculated from the averaged F and from the F calculated in (1).

The solution:

(1) The partition function is:

$$Z = \sum_{\sigma_1=\pm 1} \sum_{\sigma_2=\pm 1} \dots \sum_{\sigma_n=\pm 1} e^{-\beta(-J \sum_{i=1}^N \sigma_i \sigma_{i+1} (1 - \lambda n_i))}$$

Using the Transfer matrices formalism we get:

$$Z = \sum_{\sigma_1=\pm 1} \sum_{\sigma_2=\pm 1} \dots \sum_{\sigma_n=\pm 1} \langle \sigma_1 | P_1 | \sigma_2 \rangle \langle \sigma_2 | P_2 | \sigma_3 \rangle \dots \langle \sigma_n | P_n | \sigma_{n+1} \rangle$$

The transfer matrix for two spins without an atom between is:

$$P_i = \begin{pmatrix} e^{\beta J} & e^{-\beta J} \\ e^{-\beta J} & e^{\beta J} \end{pmatrix}$$

The transfer matrix for two spins with an atom between is:

$$P'_i = \begin{pmatrix} e^{\beta J(1-\lambda)} & e^{-\beta J(1-\lambda)} \\ e^{-\beta J(1-\lambda)} & e^{\beta J(1-\lambda)} \end{pmatrix}$$

Both matrices commute:

$$[P_i, P'_i] = 0$$

So the multiplication order does not matter. After changing to the diagonalized base we get for a single configuration:

$$\begin{aligned} Z &= Tr(P'^{N'} P^{N-N'}) \\ &= Tr\left(\begin{pmatrix} 2\cosh(\beta J)^{N-N'} & 0 \\ 0 & 2\sinh(\beta J)^{N-N'} \end{pmatrix} \begin{pmatrix} 2\cosh(\beta J(1-\lambda))^{N'} & 0 \\ 0 & 2\sinh(\beta J(1-\lambda))^{N'} \end{pmatrix} \right) \end{aligned}$$

Assuming $N \rightarrow \infty$ we can neglect the smaller eigenvalue so:

$$Z = \binom{N}{N'} 2\cosh(\beta J)^{N-N'} 2\cosh(\beta J(1-\lambda))^{N'}$$

Notice that a combinatorial factor was introduced in order to count the different configurations possible for a given N' .

The free energy is therefore:

$$F = -\tau \ln Z = -\tau \ln \left(\binom{N}{N'} \right) - \tau l(N - N') \ln(2\cosh(\beta J)) - \tau N' \ln(2\cosh(\beta J(1 - \lambda)))$$

(2) For any configuration with N' impurities the partition function is:

$$Z = 2\cosh(\beta J)^{N-N'} 2\cosh(\beta J(1-\lambda))^{N'}$$

The free energy for any configuration with N' impurities is (It doesn't depend on the specific configuration):

$$F = -\tau \ln Z = -\tau l(N - N') \ln(2\cosh(\beta J)) - \tau N' \ln(2\cosh(\beta J(1 - \lambda)))$$

The average free energy for a given N' is:

$$\langle F \rangle = \frac{\sum_{\text{configurations}} F_{\text{conf}}}{\sum_{\text{configurations}}} = \frac{\binom{N}{N'} \times -\tau l(N - N') \ln(2\cosh(\beta J)) - \tau N' \ln(2\cosh(\beta J(1 - \lambda)))}{\binom{N}{N'}}$$

$$\langle F \rangle = -\tau l(N - N') \ln(2\cosh(\beta J)) - \tau N' \ln(2\cosh(\beta J(1 - \lambda)))$$

The entropy difference between the two calculations:

$$\Delta S = -\frac{\partial F}{\partial \tau} + \frac{\partial \langle F \rangle}{\partial \tau} = \ln \left(\binom{N}{N'} \right)$$