

## E4550: Chemical potential: reaction in ideal quantum gases

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### The problem:

$N$  fermions of type  $F$  that have spin  $\frac{1}{2}$  are placed in a volume  $V$ . they can decay into a boson of type  $B$  that has spin 0 and a fermion of type  $A$  that has spin  $\frac{1}{2}$ . the reaction is  $F \rightarrow A + B$  and it has an energy gain of  $\epsilon_0$  (i.e.  $A + B$  have lower energy than  $F$ ). the masses are  $m_F, m_A, m_B$  respectively.

(1) Assuming ideal gases at temperature  $T$ , write the equations which determine the densities  $n_F, n_A, n_B$  in equilibrium.

(2) Write the equations of (1) at  $T = 0$  and describe the densities as functions of  $\epsilon_0$ . Find  $\epsilon_c$  such that for  $\epsilon_0 > \epsilon_c$  the number of F fermions vanishes.

(3) Assume that the condensation of bosons B occurs at  $T_c$  such that  $T_c \ll \frac{p_F^2}{2m_A} < \epsilon_0$  where  $p_F$  is the Fermi momentum of fermions A. Evaluate  $T_c$  and rewrite the condition  $T_c \ll \frac{p_F^2}{2m_A}$  in terms of the given parameters.

### The solution:

(1) The particles are fermions and bosons so clearly there densities are (assuming there is no degeneracy in the boson gas):

$$n_F = 2 \sum_p \frac{1}{e^{\beta(\frac{p^2}{2m_F} - \mu_F)} + 1}$$

$$n_A = 2 \sum_p \frac{1}{e^{\beta(\frac{p^2}{2m_A} - \mu_A)} + 1}$$

$$n_B = \sum_p \frac{1}{e^{\beta(\frac{p^2}{2m_B} - \mu_B)} - 1} + n_B^0$$

Where  $n_B^0$  is the bose condensation related term.

(2) At  $T = 0$  the bosons undergo a condensation meaning that

$$n_B = n_B^0$$

$$\mu_B = 0$$

The chemical potential must obey the reaction by

$$\mu_F + \epsilon_0 = \mu_A + \mu_B$$

and at  $T = 0$

$$\mu_A = \mu_F + \epsilon_0$$

in addition, the fermionic density must obey the reaction by

$$n = n_A + n_F = \text{const}$$

therefore at  $T = 0$  the fermionic density is

$$n = \frac{1}{3\pi^2\hbar^3}(2m_F\mu_F)^{\frac{3}{2}} + \frac{1}{3\pi^2\hbar^3}(2m_A\mu_A)^{\frac{3}{2}}$$

From all of the above it's clear that the condition for  $n_F = 0$  is  $\mu_F < 0$  leading to

$$0 = \mu_A - \epsilon_c \rightarrow n = \frac{1}{3\pi^2\hbar^3}(2m_A\epsilon_c)^{\frac{3}{2}}$$

The density  $\frac{n_F}{n}$  starts from the value  $\frac{n_F}{n} = \frac{m_F^{\frac{3}{2}}}{m_F^{\frac{3}{2}} + m_A^{\frac{3}{2}}}$  at  $\epsilon_0 = 0$  and decrease until it vanishes at  $\epsilon_0 = \epsilon_c$ .

$n_A = n - n_F$  and therefore the density  $\frac{n_A}{n}$  starts from the value  $\frac{n_A}{n} = \frac{m_A^{\frac{3}{2}}}{m_F^{\frac{3}{2}} + m_A^{\frac{3}{2}}}$  at  $\epsilon_0 = 0$  and increase until it gets the constant value 1 at  $\epsilon_0 = \epsilon_c$ .

(3) Under the given assumption of  $T_c \ll \frac{p_F^2}{2m_A} < \epsilon_0$  and using the fact that  $\frac{p_F^2}{2m_A} = \mu_F + \epsilon_0$  at  $T = 0$  we get that

$$\frac{p_F^2}{2m_A} = \mu_F + \epsilon_0 < \epsilon_0 \rightarrow \mu_F < 0$$

$\mu_F < 0$  leads us to  $n_F = 0$ , leading to  $n_A = n_B = n$  which eventually leads us to

$$\mu_F + \epsilon_0 = \frac{(3\pi^2\hbar^3n)^{\frac{2}{3}}}{2m_A}$$

putting that back in the condition we are given gives us

$$T_c \ll \frac{(3\pi^2\hbar^3n)^{\frac{2}{3}}}{2m_A}$$

From the other hand it's known that since  $n_B = n$

$$T_c = \frac{2\pi\hbar^2n^{\frac{2}{3}}}{m_B(2.612)^{\frac{2}{3}}}$$

substituting this to the former condition we get the condition in the given parameters

$$\frac{m_B}{m_A} \ll \frac{4\pi}{(3\pi^2 \cdot 2.612)^{\frac{2}{3}}} = 0.69$$