

Ex4220: Chemical equilibrium: $H_2[3D] \rightleftharpoons 2H[2D]$

Submitted by: Edo Hogeg

The problem:

An H_2 molecule (mass $2m_H$) decomposes into H atoms when it is absorbed upon a certain metallic surface. The H_2 and H have an energy gain ε_2 and ε_1 per atom respectively, related to each state. The binding of the H atoms is not to particular sites on the surface: the H atoms are free to move on the surface. Consider the H_2 and H as an ideal gas, and express the surface density of the H atoms as a function of the H_2 pressure.

The solution:

When the H_2 and H come to equilibrium (that is the system's free energy is minimized) their chemical potential are equal to zero. We'll start by finding the partition function.

$$Z_{1H_2} = \int \frac{d^3x d^3p}{(2\pi)^3} e^{-\beta(\frac{p^2}{2(2m)} + \varepsilon_2)} = e^{-\beta\varepsilon_2} V \left(\frac{m}{\beta\pi} \right)^{\frac{3}{2}} = e^{-\beta\varepsilon_2} \frac{2^{\frac{3}{2}} V}{\lambda_T^3} \quad (1)$$

$$Z_{1H} = \int \frac{d^2x d^2p}{(2\pi)^2} e^{-\beta(\frac{p^2}{2m} + \varepsilon_1)} = e^{-\beta\varepsilon_1} A \frac{m}{2\beta\pi} = e^{-\beta\varepsilon_1} \frac{A}{\lambda_T^2} \quad (2)$$

Now let's assume there are N H_2 molecules to begin with. With the definition of a macroscopic reaction coordinate n (See lecture notes: [5.3] The law of mass action) we know that in equilibrium there are $N - n$ H_2 molecules and $2n$ H atoms. Recalling both H_2 and H are ideal gasses we obtain:

$$Z_{H_2} = \frac{1}{(N - n)!} Z_{1H_2}^{N-n} = \frac{e^{-\beta\varepsilon_2(N-n)}}{(N - n)!} \left(V \left(\frac{m}{\beta\pi} \right)^{\frac{3}{2}} \right)^{N-n} \quad (3)$$

$$Z_H = \frac{1}{(2n)!} Z_{1H}^{2n} = \frac{e^{-\beta\varepsilon_1(2n)}}{(2n)!} \left(A \frac{m}{2\beta\pi} \right)^{2n} \quad (4)$$

$$F_{H_2} = -\frac{1}{\beta} \ln(Z_{H_2}) = -\frac{1}{\beta} \left[-\beta\varepsilon_2(N - n) - (N - n) \ln(N - n) + (N - n) + (N - n) \ln \left(V \left(\frac{m}{\beta\pi} \right)^{\frac{3}{2}} \right) \right] \quad (5)$$

$$F_H = -\frac{1}{\beta} \ln(Z_H) = -\frac{1}{\beta} \left[-\beta\varepsilon_1(2n) - 2n \ln(2n) + 2n + 2n \ln \left(A \frac{m}{2\beta\pi} \right) \right] \quad (6)$$

Note that Stirling's approximation $\ln(N!) \approx N \ln(N) - N$ was used in (5) and (6). We would like to first find the H_2 pressure, Then we'll equilibrate the chemical potential to find our desired solution.

$$P = -\frac{\partial F_{H_2}}{\partial V} = \frac{N - n}{\beta V} \quad (7)$$

$$\mu_{H_2} + \mu_H = 0 \quad (8)$$

$$\frac{\partial F_{H_2}}{\partial n} + \frac{\partial F_H}{\partial n} = 0 \quad (9)$$

$$-2\ln\left(\frac{2n}{A}\right) - 2\beta\varepsilon_1 + 2\ln\left(\frac{m}{2\pi\beta}\right) + \ln\left(\frac{N-n}{V\beta}\right) + \ln(\beta) + \beta\varepsilon_2 - \frac{3}{2}\ln\left(\frac{m}{\pi\beta}\right) = 0 \quad (10)$$

Denoting the density of the H atoms:

$$\mathcal{N} = \frac{2n}{A} \quad (11)$$

Finally the density \mathcal{N} of the H atoms is:

$$\mathcal{N}(P) = \frac{e^{-\frac{\beta}{2}(2\varepsilon_1 - \varepsilon_2)}}{2^{\frac{3}{4}}} \sqrt{\frac{\beta P}{\lambda_T}} \quad (12)$$