E3010: Heat capacity of ideal Bose gas

Submitted by: Submitted by Yigal Tirosh

The problem:

N bosons with spin 1 and mass m are in a box with a volume V. The gas is in a thermic equilibrium in a temperature.

- (1) Write an explicit expression for Tc using N, m and V.
- (2) Calculate the chemical potential, the energy and the pressure in the Boltzmann approximation $T\gg Tc$.
- (3) Calculate the chemical potential, the energy and the pressure in the regime T < Tc.
- (4) calculate C_v in the temperature $T < T_c$
- (5) What is the result for C_v in the temperatures T = Tc and in Boltzmann approximation $T \gg Tc$. (6) Show that $\frac{C_p}{C_v} = \frac{5g_{\frac{5}{2}}(z)g_{\frac{1}{2}}(z)}{3g_{\frac{3}{2}}^2(z)}$.

Why is $C_p \to \infty$ in the condensed phase?.

(7) Find γ in the adiabatic equation of state. Note that in general $\gamma \neq \frac{C_p}{C_r}$.

The solution: Note an impoved version in [ex3009].

 $(1)T = T_c$

$$T_c = \frac{1}{2m\pi} \left(\frac{N}{3V\zeta(\frac{3}{2})}\right)^{\frac{2}{3}} = \frac{1}{2m\pi} \left(\frac{N}{3V2.612}\right)^{\frac{2}{3}}.$$
 (1)

(2) $T \gg T_c$ Boltzmann approximation.

$$\frac{N}{V} = 3\frac{z}{\lambda^3} \to z = \frac{N\lambda^3}{3V}.$$
 (2)

$$E = 3\frac{3}{2}V\frac{T}{\lambda^3}z = \frac{9}{2}V\frac{T}{\lambda^3}e^{\beta\mu} = \frac{3}{2}NT.$$
 (3)

$$P = \frac{2E}{3V} = 3\frac{T}{\lambda^3}e^{\beta\mu} = \frac{NT}{V}.\tag{4}$$

(3) $T < T_c$.

$$\mu = 0. \tag{5}$$

$$E = 3V(2m)^{\frac{3}{2}} 2\pi \Gamma(\frac{5}{2})\zeta(\frac{5}{2})T^{\frac{5}{2}} = 3V(2m)^{\frac{3}{2}} 2\pi \frac{3}{4}\sqrt{\pi}1.341T^{\frac{5}{2}}.$$
 (6)

$$P = \frac{2E}{3V} = 2(2m)^{\frac{3}{2}} 2\pi \frac{3}{4} \sqrt{\pi} 1.341 T^{\frac{5}{2}}.$$
 (7)

(4) $T < T_c$

$$C_v = \frac{\partial E}{\partial T} = 3V(2m)^{\frac{3}{2}} 2\pi \Gamma(\frac{5}{2}) \zeta(\frac{5}{2}) T^{\frac{3}{2}} \frac{5}{2} = 3V(2m)^{\frac{3}{2}} 2\pi \sqrt{\pi} 1.341 T^{\frac{3}{2}} \frac{15}{8}.$$
 (8)

$$N_e = N \frac{T}{T_c}^{\frac{3}{2}} \tag{9}$$

particles in the excited states increased with T.

 $C_v \propto N_e(T)$.

 C_v increased with N_e until $N_e(T_c = T) = N$.

(5) $T = T_c$.

$$C_v = \frac{\partial E}{\partial T} = 3V(2m)^{\frac{3}{2}} 2\pi \Gamma(\frac{5}{2}) \zeta(\frac{5}{2}) T_c^{\frac{3}{2}} \frac{5}{2} = 3V(2m)^{\frac{3}{2}} 2\Pi \sqrt{\pi} 1.341 T_c^{\frac{3}{2}} \frac{15}{8} = 1.925 NK.$$
 (10)

 $C_v(T=T_c)$ Heat capacity have maximum, phase transition.

 $T > T_c$.

$$C_v = \frac{\partial}{\partial T} \left[\left(T \frac{3g_{\frac{5}{2}}(z)}{2g_{\frac{3}{2}}(z)} \right) \right]_v = \frac{15g_{\frac{5}{2}}}{4g_{\frac{3}{2}}(z)} - \frac{9g_{\frac{3}{2}}(z)}{4g_{\frac{1}{2}}(z)}. \tag{11}$$

for $T = T_c \rightarrow C_v = 1.925NK$.

 C_v grow smaller with T.

 $T \gg T_c$. Boltzmann approximation.

$$E = \frac{3}{2}NT \to C_v = \frac{3}{2}N. \tag{12}$$

(6)

$$\frac{C_p}{NK} = \frac{\partial}{\partial T} \left[\left(T \frac{5g_{\frac{5}{2}}(z)}{2g_{\frac{3}{2}}(z)} \right) \right]_p = \left[\frac{\partial z}{\partial T} \right]_p \left[\frac{\partial F(g)}{\partial z} \right]_p + F(g) = \frac{25}{4} \frac{g_{\frac{5}{2}}^2(z)g_{\frac{1}{2}}(z)}{g_{\frac{3}{2}}^3(z)} - \frac{15}{4} \frac{g_{\frac{5}{2}}(z)}{g_{\frac{3}{2}}(z)}.$$
(13)

$$F(g) = \frac{5g_{\frac{5}{2}}(z)}{2g_{\frac{3}{2}}(z)}. (14)$$

$$\left[\frac{\partial z}{\partial T}\right]_p = \frac{-5z}{2T} \frac{g_{\frac{5}{2}}(z)}{g_{\frac{3}{2}}(z)}.\tag{15}$$

 C_v From (11)

 $C_p(T < T_c) \to \infty$ all particles in the condensed phase and they compress to zero volume the system have only one phase.

(7) Adiabat of an ideal Bose gas.

For $T \leq T_c$

S - Entropy

$$\frac{S}{NK} = \frac{5}{2} \frac{V}{\lambda^3} \zeta(\frac{5}{2}) \to VT^{\frac{3}{2}} = const. \tag{17}$$

$$P = \frac{KT}{\lambda^3} \zeta(\frac{5}{2}) \to \frac{P}{T^{\frac{5}{2}}} = const. \tag{18}$$

From (17) and (18)

$$PV^{\frac{5}{3}} = const \to \gamma = \frac{5}{3}.$$
 (19)

Ideal Bose gas only for $T \gg T_c$ then $\gamma \simeq \frac{5}{3}$.