

E901: Scattering of particles by a simple cubic lattice

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Problem:

Particles are scattered from 3D cubic lattice, with the characteristic distance between atoms a . The lattice contain $(2N + 1)^3$ atoms, location of the i -atom:

$$\vec{a}_n = a(n_x\vec{e}_x + n_y\vec{e}_y + n_z\vec{e}_z) \quad -N \leq n_x, n_y, n_z \leq N$$

the atoms are identical and their potentials $U(r)$ are identical as well.

- a) find the differential cross section for the single atom in a first order Born Approximation.
- b) find the scattering amplitude $f(\Omega)$ and differential cross section $\frac{d\sigma}{d\Omega}$ for finite N
- c) find the differential cross section $\frac{d\sigma}{d\Omega}$ for infinite N

Solution:

The total potential could be written:

$$V(r) = \sum_{n_x n_y n_z} U(r - an) \quad (1)$$

We define:

$$\vec{q} = \vec{k}_\Omega - \vec{k}_0$$

$$\tilde{V}(q) = F.T.[V(r)]$$

$$\tilde{U}(q) = F.T.[U(r)]$$

Born approximation for the scattering amplitude:

$$f(\Omega) = -\frac{m}{2\pi\hbar^2} \int d^3r e^{iqr} V(r) = -\frac{m}{2\pi\hbar^2} \tilde{V}(q) \quad (2)$$

We may look at the total potential as a convolution of a single atom potential and an array of $2N+1$ delta functions, which we call comb, then we have $\text{comb}(x)$, $\text{comb}(y)$ and $\text{comb}(z)$:

$$V(r) = U(r) (*) \text{comb}(x) \text{comb}(y) \text{comb}(z) \quad (3)$$

According to the properties of F.T.:

$$\tilde{V}(q) = F.T.[V(r)] = F.T.[U(r)] \cdot F.T.[\text{comb}(x)] \cdot F.T.[\text{comb}(y)] \cdot F.T.[\text{comb}(z)] \quad (4)$$

a.

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2 = \frac{m^2}{4\pi^2\hbar^4} \tilde{U}^2(q) \quad (5)$$

b.

We find the F.T.[comb(x)] for finite N and then take it to infinity (see the Appendix), and the results are:

the scattering amplitude for a finite N (see Appendix 1):

$$f(\Omega) = -\frac{m}{2\pi\hbar^2}\tilde{U}(q)\frac{\sin q_x a(N+1/2)}{\sin \frac{q_x a}{2}}\frac{\sin q_y a(N+1/2)}{\sin \frac{q_y a}{2}}\frac{\sin q_z a(N+1/2)}{\sin \frac{q_z a}{2}}, \quad (6)$$

the differential cross section for a finite N:

$$\frac{d\sigma}{d\Omega} = |f(\Omega)|^2 = \frac{m^2}{4\pi^2\hbar^4}\tilde{U}^2(q)\frac{\sin^2 q_x a(N+1/2)}{\sin^2 \frac{q_x a}{2}}\frac{\sin^2 q_y a(N+1/2)}{\sin^2 \frac{q_y a}{2}}\frac{\sin^2 q_z a(N+1/2)}{\sin^2 \frac{q_z a}{2}}, \quad (7)$$

c.

Here we check the limit $N \rightarrow \infty$ and obtain the next result (see Appendix 2):

$$\frac{d\sigma}{d\Omega} = \frac{m^2}{4\pi^2\hbar^4}\tilde{U}^2(q)\frac{(2\pi)^6}{a^6}\sum_{n_x}\delta(q_x - \frac{2\pi}{a}n_x)\sum_{n_y}\delta(q_y - \frac{2\pi}{a}n_y)\sum_{n_z}\delta(q_z - \frac{2\pi}{a}n_z) \quad (8)$$

From here we obtain the Bragg condition for scattering:

$$q = \frac{2\pi n}{a} \text{ or } 2k \sin \frac{\theta}{2} = \frac{2\pi}{a}n$$

$$\text{Definition: } \lambda = \frac{2\pi}{k}, \alpha = \frac{\theta}{2}$$

we obtain the known formula:

$$2a \sin \alpha = n\lambda$$

The indexes n_x, n_y, n_z are called the Miller's Indices.

Appendix 1

Calculation of F.T.[comb(x)]:

$$comb(x) = \sum_{n_x=-N}^N \delta(x - na)$$

$$Comb(q) = F.T.[comb(x)] = e^{-iqaN} + e^{-iqa(N+1)} + .. + 1 + .. + e^{iqa(N-1)} + e^{iqaN}$$

we use the known formula for the sum of the geometric progression $S_n = b_1 \frac{q^N - 1}{q - 1}$:

$$\begin{aligned} &= |\alpha = qa| = e^{+i\alpha N} \left(\frac{e^{-i\alpha(2N+1)} - 1}{e^{-i\alpha} - 1} \right) = \frac{e^{-i\alpha(N+1)} - e^{+i\alpha N}}{e^{-i\alpha} - 1} = \frac{e^{-i\alpha(N+1)/2} - e^{+i\alpha(N+1)/2}}{e^{-i\alpha/2} - e^{+i\alpha/2}} = \\ &= \frac{\sin \alpha(N + \frac{1}{2})}{\sin \frac{\alpha}{2}} \end{aligned}$$

It is easy to show (basic trigo) that under transformation $\alpha \rightarrow \alpha - 2\pi\nu$ (ν is an integer)
The expression $\frac{\sin^2 \alpha(N+1/2)}{\sin^2 \alpha/2}$ does not change.

Appendix 2

We want to analyze the expression

$$\frac{\sin^2 \alpha(N + 1/2)}{\sin^2 \frac{\alpha}{2}}$$

This expression has the poles at $\alpha = 2\pi n$ and the limit at $\alpha \rightarrow 2\pi n$ is
(according to L'Hopitals Rule):

$$\begin{aligned} \lim_{\alpha \rightarrow 2\pi n} \frac{\sin^2 \alpha(N + 1/2)}{\sin^2 \frac{\alpha}{2}} &= \lim_{\alpha \rightarrow 2\pi n} \frac{(N + \frac{1}{2})2 \sin \alpha(N + 1/2) \cos \alpha(N + \frac{1}{2})}{\sin \frac{\alpha}{2} \cos \frac{\alpha}{2}} = \\ \lim_{\alpha \rightarrow 2\pi n} \frac{(N + \frac{1}{2}) \sin 2\alpha(N + \frac{1}{2})}{\frac{1}{2} \sin \alpha} &= \lim_{\alpha \rightarrow 2\pi n} \frac{4(N + \frac{1}{2})^2 \cos 2\alpha(N + \frac{1}{2})}{\cos \alpha} = 4(N + \frac{1}{2})^2 = (2N + 1)^2 \end{aligned}$$

If we have a large number of atoms ($N \rightarrow \infty$) and making the use of the fact

$$\frac{1}{2\pi N} \int_{-\infty}^{\infty} \frac{\sin^2 Nx}{x^2} = 1$$

We come to the conclusion that under this conditions ($\alpha \rightarrow 2\pi n, N \rightarrow \infty$) we get an array of delta functions.

Then

$$F.T.[comb(x)] = Comb(q) = \frac{2\pi}{a} \sum_n \delta(q - \frac{2\pi}{a}n)$$

As it was expected from the direct F.T. of infinite array of delta functions