

E6514: WKB Transmission of Harmonic and Morse Barriers

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The problem:

The exact "under the barrier" transmission coefficients for the repulsive 1D Harmonic $-(1/2)Kx^2$ and Morse $(U_0/\cosh^2(\alpha x))$ potentials can be found in [Landau and Lifshitz]¹. Obtain the zero order WKB approximation coefficients for these potentials and compare them to the exact results.

The solution:

The transmission coefficient for a potential barrier entirely above the energy level under the zero order WKB approximation is given by,

$$T = \exp\left(-\frac{2}{\hbar} \int_{x_1}^{x_2} \sqrt{2m(V(x) - E)} dx\right)$$

The integration here occurs between the classical turning points. This approximation holds for sufficiently small values of T (That is, well below the peak in this problem) and sufficiently long wavelengths.

(1) The transmission coefficient for the 1D Repulsive Harmonic Potential given in [L.L] at §50 Problem 4, under the condition that $E_0/\hbar \gg \sqrt{K/m}$ where $E_0 = |E|$, becomes

$$T = \frac{1}{1 + e^{2\pi(E_0/\hbar)\sqrt{m/K}}} \approx e^{-2\pi(E_0/\hbar)\sqrt{m/K}}$$

The same result is obtained using WKB by evaluating the following integral between the two classical turning points $(x_{1,2} = \pm\sqrt{2E/K})$ using a $\sin \alpha = \sqrt{K/2E_0}x$ substitution,

$$\phi = \int_{x_1}^{x_2} \sqrt{2m\left(E_0 - \frac{1}{2}Kx^2\right)} dx = 2E_0\sqrt{\frac{m}{K}} \int_0^\pi \sin^2 \alpha d\alpha = \pi E_0\sqrt{\frac{m}{K}}$$

and using the result in $T = e^{-2\phi/\hbar}$.

(2) The transmission coefficient for the 1D Repulsive Morse Potential given in [L.L] at §25 Problem 4, under the condition that $U_0 \gg E \gg \hbar^2\alpha^2/2m$, becomes

$$T = \left(1 + \left(\frac{\cosh\left(\frac{\pi}{2}\sqrt{8mU_0/\hbar^2\alpha^2 - 1}\right)}{\sinh\left(\pi\sqrt{2mE}/\hbar\alpha\right)}\right)^2\right)^{-1} \approx -e^{2\pi(\sqrt{2mU_0} - \sqrt{2mE})/\hbar\alpha}$$

This result can also be obtained using WKB by evaluating the following integral between the two classical turning points (Which satisfy $C^2 \equiv U_0/E - 1 = \sinh^2(\alpha x)$) using a $t = \sinh(\alpha x)$ substitution leading to an unpleasant contour integral with a known result,

$$\phi = \int_{x_1}^{x_2} \sqrt{2m\left(\frac{U_0}{\cosh^2 \alpha x} - E\right)} dx = \frac{\sqrt{2mE}}{\alpha} \int_{-C}^C \frac{\sqrt{C^2 - t^2}}{1 + t^2} dt = \frac{\sqrt{2mE}}{\alpha} \pi \left(\sqrt{C^2 + 1} - 1\right)$$

and using this result in $T = e^{-2\phi/\hbar}$.

¹L. D. Landau, E.M. Lifshitz "Quantum Mechanics: Non-Relativistic Theory. Vol. 3"