

## Ex5662: Landau levels in Graphene

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### The problem:

The effective Hamiltonian of an electron in a 2-D layer of Graphene is  $H = v_0 \boldsymbol{\sigma}(\mathbf{p} - e\mathbf{A})$  while  $\boldsymbol{\sigma} = (\sigma_x, \sigma_y, 0)$  and  $\mathbf{A}$  is the vector potential for a perpendicular magnetic field  $\mathbf{B}$  within the Landau gauge.

Notice that the standard basis for an electron is  $|x, y, m\rangle$ ,  $m = \downarrow, \uparrow$ .

(1) In the absence of a magnetic field, for a given momentum  $\mathbf{p} = (p_x, p_y)$ , what are the eigenenergies of the particle?

(2) Within the Landau gauge  $\hat{Y} = -(\frac{1}{eB})\hat{p}_x$  is a constant of motion. Write down the Hamiltonian  $H^{(Y)} = H_{m,m'}(p_y, y - Y)$  after the variable separation. Note that the linear combination of canonical coordinates  $a = \frac{(sQ + \frac{iP}{s})}{\sqrt{2}}$  is a ladder operator, and write down the Hamiltonian derived in the alternative notation  $H_{m,m'}(a, a^\dagger)$ .

(3) Define the operator  $C = (H^{(Y)})^2$ .

Calculate the eigenvalues  $\lambda_{n=0,1,2,\dots}$  of  $C$ ,

notice, the eigenvalues are degenerated except  $n = 0$ .

(4) Since  $C$  is a constant of motion it is possible to use variable separation once more.

Write down the  $2 \times 2$  matrix representing  $H^{(Y,n)}$ .

(5) What are the energy levels  $E_{Y,n,\pm}$  for  $n > 0$ ?

(6) Find the eigenstates and write them in the standard basis.

Notice that the quantum state in the standard basis is represented by  $\Psi \mapsto (\psi_\uparrow(x, y), \psi_\downarrow(x, y))$ .

You may use the notation  $\varphi^n()$  for the eigenfunctions of a 1-D harmonic oscillator.

### The solution:

(1) The Hamiltonian without a magnetic field is:

$$\begin{aligned} H &= v_0 \boldsymbol{\sigma} \cdot \mathbf{p} = v_0 (\sigma_x p_x + \sigma_y p_y) \\ &= v_0 \begin{pmatrix} 0 & p_x - ip_y \\ p_x + ip_y & 0 \end{pmatrix} = v_0 |\mathbf{p}| \begin{pmatrix} 0 & e^{-i\phi} \\ e^{i\phi} & 0 \end{pmatrix} \end{aligned}$$

where  $\mathbf{p} = (p_x, p_y)$  is given and  $tg(\phi) = \frac{p_y}{p_x}$ .

Diagonalization of this Hamiltonian provides the following eigenenergies:  $E_\pm = \pm v_0 |\mathbf{p}|$ .

(2) The vector potential within the Landau gauge is:

$$\mathbf{A} = B(-y, 0, 0)$$

Therefore the Hamiltonian is:

$$H = v_0 [\sigma_x (p_x + eB_y) + \sigma_y p_y]$$

plugging the definition of  $Y$  into  $H$  yields:

$$H = v_0 [eB \sigma_x (y - Y) + \sigma_y p_y]$$

One should notice that  $[Y, H] = 0$ , thus the Hamiltonian is separable into blocks of the form:

$$H^{(Y)} = v_0 [eB \sigma_x (y - Y) + \sigma_y p_y] = v_0 \sqrt{eB} [\sqrt{eB} \sigma_x (y - Y) + \frac{1}{\sqrt{eB}} \sigma_y p_y]$$

We define:

$$Q = y - Y, P = p_y, s = \sqrt{eB}$$

So:

$$H^{(Y)}(Q, P) = v_0[\sigma_x \sqrt{eB} Q + \frac{1}{\sqrt{eB}} \sigma_y P]$$

Using the ladder operator given:

$$a = \frac{1}{\sqrt{2}}(Q\sqrt{eB} + \frac{i}{\sqrt{eB}}P)$$

$$\Downarrow$$

$$Q = \frac{1}{\sqrt{2eB}}(a^\dagger + a), P = \frac{\sqrt{eB}}{i\sqrt{2}}(a - a^\dagger)$$

plugging into the Hamiltonian:

$$H^{(Y)}(a, a^\dagger) = v_0 \sqrt{2eB} [a(\sigma_x - i\sigma_y) + a^\dagger(\sigma_x + i\sigma_y)]$$

$$H^{(Y)} \mapsto v_0 \sqrt{2eB} \begin{pmatrix} 0 & a^\dagger \\ a & 0 \end{pmatrix}$$

(3)

$$[a, a^\dagger] = 1 \Rightarrow aa^\dagger = 1 + a^\dagger a$$

$$C = (H^{(Y)})^2 \mapsto 2eBv_0^2 \begin{pmatrix} a^\dagger a & 0 \\ 0 & aa^\dagger \end{pmatrix} = 2eBv_0^2 \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a + 1 \end{pmatrix}$$

The eigenstates of  $a^\dagger a$  satisfy: (1D harmonic oscillator)

$$a^\dagger a |n\rangle = n |n\rangle, (n = 0, 1, 2 \dots)$$

The eigenvectors of  $C$  are:  $\begin{pmatrix} |n\rangle \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ |n-1\rangle \end{pmatrix}$

$$C^{(n)} \mapsto 2eBv_0^2 \begin{pmatrix} n & 0 \\ 0 & n \end{pmatrix} \Rightarrow \lambda_n = 2eBv_0^2 n$$

(4)  $[C, H^{(Y)}] = 0$ , therefore  $H^{(Y)}$  is separable into blocks.

$$H^{(Y)} \begin{pmatrix} |n\rangle \\ 0 \end{pmatrix} = v_0 \sqrt{2eB} \begin{pmatrix} 0 \\ a |n\rangle \end{pmatrix} = v_0 \sqrt{2eBn} \begin{pmatrix} 0 \\ |n-1\rangle \end{pmatrix}$$

$$H^{(Y)} \begin{pmatrix} 0 \\ |n-1\rangle \end{pmatrix} = v_0 \sqrt{2eB} \begin{pmatrix} a^\dagger |n-1\rangle \\ 0 \end{pmatrix} = v_0 \sqrt{2eBn} \begin{pmatrix} |n\rangle \\ 0 \end{pmatrix}$$

$$\Downarrow$$

$$H^{(Y,n)} \mapsto v_0 \sqrt{2eB} \begin{pmatrix} 0 & \sqrt{n} \\ \sqrt{n} & 0 \end{pmatrix}$$

(5)

$$H^{(Y,n)} = v_0 \sqrt{2eBn} \sigma_x \Rightarrow E_{Y,n,\pm} = \pm v_0 \sqrt{2eBn}, (n > 0)$$

(6) The eigenstates of  $H^{(Y)}$  are (up to a normalization factor):

$$\begin{pmatrix} |n\rangle \\ \pm |n-1\rangle \end{pmatrix}$$

the representation of these eigenstates in the standard basis is:

$$\begin{pmatrix} \varphi^n(y) \\ \pm \varphi^{n-1}(y) \end{pmatrix}$$

remembering that  $Y$  shares the same eigenstates with  $p_x$  yields the following eigenstates:

$$e^{-ieBYx} \begin{pmatrix} \varphi^n(y) \\ \pm \varphi^{n-1}(y) \end{pmatrix}$$