

Ex5662: Landau Levels in Grafene

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The problem:

The effective Hemiltonian of an electron in a 2-D layer of Grafene is $H = v_0\sigma(p - eA)$ when $\sigma = (\sigma_x, \sigma_y, \sigma_z)$ and A is the vector potential for a pependicular magnetic field B in Landau gauge. Be advised that the standard base for representing the electron is $|x, y, m\rangle$, when $m = \downarrow, \uparrow$.

- (1) In the lack of a magnetic field, and for a given momentum $p = (p_x, p_y)$ what are the eigen energies of the particle?
- (2) In Landau gauge $\hat{Y} = -(1/(eB))\hat{p}_x$ is a constant of motion. Write the Hemiltonian $H^Y = H_{m,m}(p_y, y - Y)$ of after the seperation of variables. Please note that a linear combination of canonical coordinates $a = (sQ + iP/s)/\sqrt{2}$ is a ladder operator, and write the Hemiltonian in the alternative way $H_{m,m}(a, a^\dagger)$
- (3) Define the operator $C = (H^Y)^2$ and calculate the eigenvalues $\lambda_n = 0, 1, 2, \dots$ of C . Please note that all the eigenvalue except $n=0$ are degenerated.
- (4) Since C is a constant of motion it is possible to perform a second seperation of variables. Write the 2×2 matrix representing $H^{Y,n}$.
- (5) Find the energy levels $E_{Y,n,\pm}$ for $n > 0$
- (6) Write the matrix form of the operator $\hat{I}(t)$ and represent it as a sum of Pauli Matrices. Find the eigenstates and write them in the standard base. Be advised that the quantum state in the standard base is represented by $\bar{\Theta} \rightarrow (\Theta_\uparrow(x, y), \Psi_\downarrow(x, y))$. It is possible to make use of $\varphi^n()$ for the eigen functions of a 1-D harmonic oscillator.

The solution:

Preface: $H = V_0\sigma \cdot (p - e\bar{A})$, $\bar{B} = B\hat{z} \implies \{Landau - Gauge\} \implies \bar{A} = (-By, 0, 0)$

$$\begin{aligned} \sigma &= (\sigma_x, \sigma_y), \quad p = (p_x, p_y) \implies \sigma \cdot p = \sigma_x p_x + \sigma_y p_y, \quad \sigma \cdot e\bar{A} = -eBy\sigma_x \\ \implies H &= v_0[\sigma_x(p_x + eBy) + \sigma_y p_y] = v_0 \begin{pmatrix} 0 & (p_x + eBy - ip_y) \\ (p_x + eBy + ip_y) & 0 \end{pmatrix} \\ \implies E &= \pm v_0(p_x^2 + 2p_x eBy + (eBy)^2 + p_y^2)^{1/2} = \pm v_0(\bar{p}^2 + 2p_x eBy + (eBy)^2)^{1/2} \end{aligned}$$

$$(1) \text{ for } B = 0 \implies E = \pm v_0 |\bar{p}|$$

$$(2) \quad H = V_0\sigma(\bar{p} + eBy\hat{x}), \text{ Let us define } \hat{Y} = -\hat{p}_x/(eB) \text{ so that } [H, Y] = 0 \\ \implies H = v_0[\sigma_x eB(-Y + y) + \sigma_y p_y] = v_0 eB \begin{pmatrix} 0 & y - Y - ip_y/eB \\ y - Y + ip_y/eB & 0 \end{pmatrix}$$

Let us define Q and P as canonical conjugates in the following manner:

$$Q = y - Y, \quad P = p_y$$

and so we can define ladder operators a and a^\dagger in the following way:

$$a = (sQ + iP/s)/\sqrt{2}$$

$$a^\dagger = (sQ - iP/s)/\sqrt{2}$$

when $s = \sqrt{2eB}$ and the Hemiltonian will take the form:

$$H = v_0 \sqrt{2eB} \begin{pmatrix} 0 & a^\dagger \\ a & 0 \end{pmatrix}$$

$$(3) \quad [a, a^\dagger] = aa^\dagger - a^\dagger a = 1 \implies aa^\dagger = 1 + a^\dagger a$$

$$C = (H^Y)^2 = V_0^2 \cdot 2eB \begin{pmatrix} a^\dagger a & 0 \\ 0 & aa^\dagger \end{pmatrix} = V_0^2 \cdot 2eB \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a + 1 \end{pmatrix}$$

Let us define: $N = a^\dagger a$ so that $N |n\rangle = n |n\rangle \Rightarrow$

$$C = V_0^2 \cdot 2eB \begin{pmatrix} n & 0 \\ 0 & n+1 \end{pmatrix}$$

$$u_n = \begin{pmatrix} n \\ 0 \end{pmatrix}, u_{n+1} = \begin{pmatrix} 0 \\ n+1 \end{pmatrix}, \lambda_1 = n, \lambda_2 = n+1$$

$E_n^2 = V_0^2 \cdot 2eBn$, $\lambda_n = n = 0, 1, 2, \dots$ and E_n is of degeneracy 2 for every n but $n = 0$

$$(4) H^Y = V_0 \cdot \sqrt{2eB} \begin{pmatrix} 0 & a \\ a^\dagger & 0 \end{pmatrix}, C = (H^Y)^2 = V_0^2 \cdot 2eB \begin{pmatrix} a^\dagger a & 0 \\ 0 & a^\dagger a + 1 \end{pmatrix} \text{ so } [H, C] = 0$$

so C and H share the same eigen states;

$$u_1 = \begin{pmatrix} |n\rangle \\ 0 \end{pmatrix}, u_2 = \begin{pmatrix} 0 \\ |n+1\rangle \end{pmatrix}$$

we can rearrange the order of the base elements so that

$$C = (H^Y)^2 = V_0^2 \cdot 2eB \begin{pmatrix} n & 0 \\ 0 & n+1-1 \end{pmatrix} = V_0^2 \cdot 2eB \begin{pmatrix} n & 0 \\ 0 & n \end{pmatrix}$$

$$u_1 = \begin{pmatrix} |n\rangle \\ 0 \end{pmatrix}, u_2 = \begin{pmatrix} 0 \\ |n\rangle \end{pmatrix}$$

$$C = (H^Y)^2 \Rightarrow H^Y = \sqrt{C} = v_0 \sqrt{2eB} \begin{pmatrix} 0 & \sqrt{n} \\ \sqrt{n} & 0 \end{pmatrix}$$

we can rearrange the order of the base elements so that

$$\sqrt{C} = H^{Y,n} = v_0 \sqrt{2eB} \begin{pmatrix} 0 & \sqrt{n} \\ \sqrt{n} & 0 \end{pmatrix} \Rightarrow$$

$$(5) H^{Y,n} = V_0 \cdot \sqrt{2eBn} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, E_{Y,n} = \pm V_0 \cdot \sqrt{2eBn}$$

(6) In the first form of the Y blocks, the eigen vectors of the matrix:

$$H = v_0 \sqrt{2eB} \begin{pmatrix} 0 & a^\dagger \\ a & 0 \end{pmatrix} \text{ is } u_n^+ = \begin{pmatrix} |n\rangle \\ |n+1\rangle \end{pmatrix}, u_n^- = \begin{pmatrix} |n\rangle \\ -|n+1\rangle \end{pmatrix}$$

Using Q and P we see that the given Hemiltonian is analogous to one of a harmonic oscilator, therefore we can rewrite the eigenvectors in terms of the eigen functions of a harmonic oscilator;

$$u_n^+ = \begin{pmatrix} |\varphi^n(y)\rangle \\ |\varphi^{n-1}(y)\rangle \end{pmatrix}, u_n^- = \begin{pmatrix} |\varphi^n(y)\rangle \\ -|\varphi^{n-1}(y)\rangle \end{pmatrix}$$

Though since the oscilator is shifted we need to add a phase factor of the form e^{-ieBYx} so the eigen states of the given Hemiltonian will be:

$$\overline{u_n^+} = \begin{pmatrix} |\varphi^n(y)\rangle \\ |\varphi^{n-1}(y)\rangle \end{pmatrix} e^{-ieBYx}, \overline{u_n^-} = \begin{pmatrix} |\varphi^n(y)\rangle \\ -|\varphi^{n-1}(y)\rangle \end{pmatrix} e^{-ieBYx}$$

* if we consider the eigen functions to be normlized, another normlizing factor should be added