# E1344: Oscillations between a site and a ring Submitted by: Asaf Barak, Jeremy Gartner 

## The problem:

A given ring with length $L=1$ has $N$ sites, which have equal potential $(V=0)$. The hopping amplitude of a particle per time unit between neighbouring sites is $c$. Another site is added at the center of the ring. The relation energy of the particle in the central site is $\varepsilon_{0}$. The hopping amplitude per time unit from the central site to any of the other sites along the ring is $c_{0}$. The system is then placed in a magnetic field, so that the total flux through the ring is $\Phi$. The charge of the particle is e. From the following analysis, it is derived that the particle will perform oscillations between the site and the ring. This is a generalization of oscillations in a system of two sites.
(1) Write the Hamiltonian $H_{\text {ring }}(p)$ of a particle in the ring described above (without the central site).
(2) Write the eigenenergies $E_{n}$ of the particle in the ring (without the central site).
(3) Calculate the coupling $\left\langle E_{n}\right| H\left|\varepsilon_{0}\right\rangle$ of the central site to the states of the ring.
(4) What is the oscillation frequency $\Omega$ of the particle?
(5) For a given magnetic flux - What should be the coupled energy $\varepsilon_{0}$ in order to get full oscillations?
(6) For oscillation $\Omega(\Phi)$ there is periodical dependency in the magnetic flux. What is the period, based on Aharonov-Bohm?

Answer the questions, using given data only.

## The solution:

(1) From the lecture notes [9.1], we know that:

$$
\widehat{H}=c e^{-i a(\widehat{p}-A)}+c e^{i a(\widehat{p}-A)}+\text { const }
$$

Because of gauge consideration, we can determine that const $=0$.
In addition, we know that $a=\frac{L}{N}=\frac{1}{N}$, and $A=\frac{\mathrm{e} \Phi}{L}$. Therefore:

$$
\widehat{H}=c e^{-i \frac{1}{N}(\widehat{p}-\mathrm{e} \Phi)}+c e^{i \frac{1}{N}(\widehat{p}-\mathrm{e} \Phi)}=2 c \cdot \cos \left(\frac{1}{N}(\widehat{p}-\mathrm{e} \Phi)\right)
$$

(2) Considering the hamiltonian, the eigenstates are the momentum states and the eigenvalues are as follows:

$$
E_{n}=2 c \cdot \cos \left(\frac{1}{N}(2 \pi n-\mathrm{e} \Phi)\right)
$$

(3) Due to the fact that $\widehat{H}=\widehat{H}(\hat{p})$, the eigenstates $\left|E_{n}\right\rangle=\left|k_{n}\right\rangle$. Hence,

$$
\begin{aligned}
\left\langle E_{n}\right| H\left|\varepsilon_{0}\right\rangle & =\left\langle k_{n}\right| H\left|\varepsilon_{0}\right\rangle \\
& =\frac{1}{\sqrt{N}} \sum_{x=1}^{N} e^{i k_{n} x}\langle x| H\left|\varepsilon_{0}\right\rangle \\
& =\frac{1}{\sqrt{N}} \sum_{x=1}^{N} e^{i k_{n} x} \cdot c_{0} \\
& =\frac{1}{\sqrt{N}} N \delta_{n, 0} c_{0} \\
& =\sqrt{N} c_{0} \delta_{n, 0}
\end{aligned}
$$

(4) Due to the coupling, there will be oscillations between the central site to the ring with 0 momentum:

$$
\left(\begin{array}{cc|cccc}
\varepsilon_{0} & c_{0} \sqrt{N} & & & & \\
c_{0} \sqrt{N} & E_{0} & & & & \\
\hline & & E_{1} & & & \\
& & & E_{2} & & \\
& & & & \ddots & \\
& & & & & E_{n}
\end{array}\right)
$$

Because of gauge consideration, we can determine for the $2 \times 2$ subspace that:

$$
\left(\begin{array}{cc}
\frac{\varepsilon_{0}-E_{0}}{2} & c_{0} \sqrt{N} \\
c_{0} \sqrt{N} & -\frac{\varepsilon_{0}-E_{0}}{2}
\end{array}\right)
$$

From the lecture notes [36.1], we learn that the evolution of any system whose states form a dim $=2$ Hilbert space can always be described by using a precession picture. Therefore, the subspace can be rewritten as:

$$
\frac{\varepsilon_{0}-E_{0}}{2} \cdot \sigma_{z}+c_{0} \sqrt{N} \cdot \sigma_{x}=\vec{\Omega} \cdot \vec{S}
$$

Where $\vec{\Omega}=\left(2 c_{0} \sqrt{N}, 0, \varepsilon_{0}-E_{0}\right)$, and $E_{0}=2 c \cdot \cos \left(\frac{e \Phi}{N}\right)$.
Therefore:

$$
\Omega=\sqrt{\left(\varepsilon_{0}-2 c \cdot \cos \left(\frac{\mathrm{e} \Phi}{N}\right)\right)^{2}+4 c_{0}^{2} N}
$$

(5) Using the Rabi formula:

$$
P(t)=1-\sin ^{2}\left(\theta_{0}\right) \sin ^{2}\left(\frac{\Omega t}{2}\right)
$$

Where $\theta_{0}$ is the angle between the $\widehat{z}$ axis and the rotation axis $(\vec{\Omega})$.
The condition to get a full oscillation is $\theta_{0}=\frac{\pi}{2}$, and in our case $\theta_{0}$ is:

$$
\theta_{0}=\arctan \left(\frac{2 c_{0} \sqrt{N}}{\varepsilon_{0}-2 c \cdot \cos \left(\frac{\mathrm{e} \Phi}{N}\right)}\right)
$$

Therefore:

$$
\varepsilon_{0}=2 c \cdot \cos \left(\frac{\mathrm{e} \Phi}{N}\right)
$$

(6) The oscillation $\Omega(\Phi)$ has a period depending on the the cosine function, as we can see in the equation from section (4).
Therefore, the necessary condition is:

$$
\frac{\mathrm{e} \Phi}{N}=2 \pi k, \quad k \in \mathbb{Z} \quad \Longrightarrow \quad \Phi=\frac{2 \pi N}{\mathrm{e}} k
$$

and the condition is:

$$
\Phi \longmapsto \Phi+\frac{2 \pi N}{\mathrm{e}} k
$$

