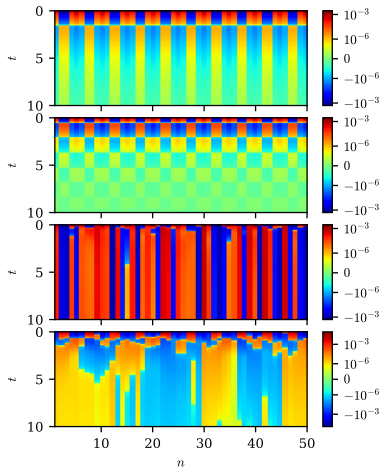
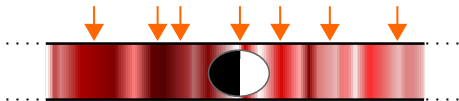


Localization due to topological stochastic disorder in active networks

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[1] DS, Dganit Meidan and Doron Cohen (Phys. Rev. E 98, 012107)

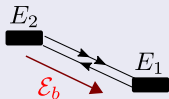
Active networks

Rate equation: $\dot{p} = Wp$

$p(t) = p^{\text{NESS}} + \sum_{\lambda} \psi_{\lambda} e^{-\lambda t}$

$W\psi_{\lambda} = -\lambda\psi_{\lambda}$

Two sites

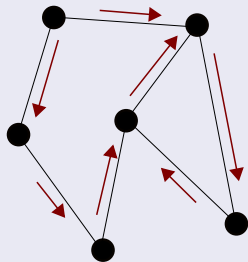


b = bond

Stochastic field $\mathcal{E}_b = \ln \left(\frac{w_{\vec{b}}}{w_{\overleftarrow{b}}} \right)$

Boltzmann $\mathcal{E}_b = \frac{E_2 - E_1}{T_b}$

Network

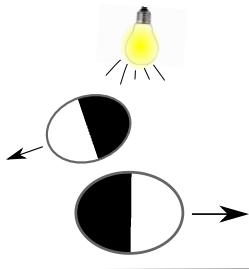
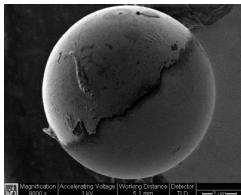


- **Affinity** $\equiv \oint \mathcal{E} dl$
- All **affinities** are 0 $\Leftrightarrow \mathcal{E}$ conservative $\Rightarrow \lambda$ are real
(Example: detailed balance)

Active Network: Non-zero **affinities** $\Rightarrow \lambda$ might be complex, under-damped relaxation

Janus particles

- Spherical-like “nano-particles” ($100\text{nm} - 10\mu\text{m}$), coated at each of their two hemispheres with different materials
- Placed in solution - diffusion
- Due to asymmetry - can be made to self propel (“active particle”)
- Sample mechanism: self-propelled when radiated with light by thermophoresis [1]



[1] Jiang, Hong-Ren, Natsuhiko Yoshinaga, and Masaki Sano. “Active motion of a Janus particle by self-thermophoresis in a defocused laser beam.” PRL 105.26 (2010): 268302.

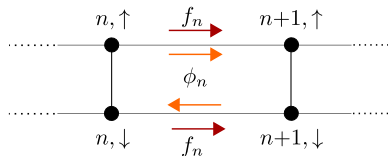
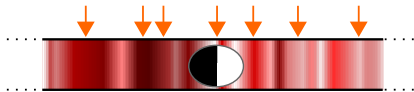
[3] Ben Yellen, Duke University

Janus particle - minimal model

- Rate equation - quasi 1D network
- Janus 1D $\Leftrightarrow |n, s\rangle$
Position: $n = 1, 2, \dots, N$
Polarization: $s = \uparrow, \downarrow$

Stochastic field on bond (n, s) :

- Drift: $f_n = \bar{f} + [-\sigma_f, \sigma_f]$
Conservative Stochastic Disorder (CSD)
- Propulsion: $\phi_n = \bar{\phi} + [-\sigma_\phi, \sigma_\phi]$
Topological Stochastic Disorder (TSD)



$$\mathcal{E}_{n,\uparrow} = f_n + \phi_n$$

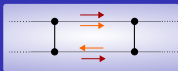
$$\mathcal{E}_{n,\downarrow} = f_n - \phi_n$$

Affinity in unit cell = $2\phi_n$

$$\mathbf{W} = \sigma_{\mathbf{w}} + \mathbf{W}_{\text{hop}} - \sum_{n,s} |n, s\rangle \gamma_{n,s} \langle n, s|$$

$$\mathbf{W}_{\text{hop}} = \sum_{n,s} |n+1, s\rangle \langle n, s| e^{\frac{\mathcal{E}_{n,s}}{2}} + |n, s\rangle \langle n+1, s| e^{-\frac{\mathcal{E}_{n,s}}{2}}$$

$$\gamma_{n,s} = 1 + e^{\mathcal{E}_{n,s}/2} + e^{-\mathcal{E}_{n-1,s}/2}$$



Clean system



Drift = 0, increasing propulsion ($\bar{\phi}$)

- No propulsion ($\bar{\phi} = 0$):

$$\lambda_{k,+} = 2 - 2 \cos(k) \quad k = 2\pi n/N$$

$$\lambda_{k,-} = 4 - 2 \cos(k)$$

Adding propulsion ($\bar{\phi}$)

- Bloch (two bands): $|k, s\rangle$

$$\mathbf{W}^{(k)} = b\sigma_x - ia\sigma_z + c\mathbf{1}$$

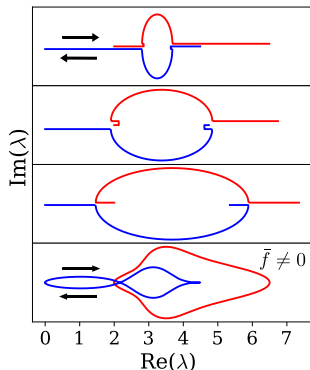
- PT symmetry breaking $a(k, \bar{\phi}) > b$

- Spectrum is complex for $\bar{\phi} > \phi_c$
($\phi_c \approx 0.96$)

- Eigenstates become polarized

$$|k, \pm\rangle = \sum_n e^{ikn} (|n, \uparrow\rangle \pm e^{\pm i\varphi} |n, \downarrow\rangle)$$

- Spectrum: $\mathbf{W}\psi = -\lambda\psi$
- Increasing propulsion ($\bar{\phi}$)



$$a = \left[2 \sinh\left(\frac{\bar{\phi}}{2}\right) \right] \sin(k) \quad c = \left[2 \cosh\left(\frac{\bar{\phi}}{2}\right) \right] \cos(k) - \left[1 + 2 \cosh\left(\frac{\bar{\phi}}{2}\right) \right]$$

$$\lambda_{k,\pm} = -[c \pm \sqrt{b^2 - a^2}] \quad b = 1$$

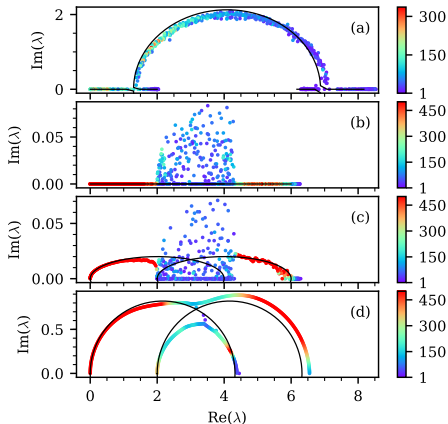
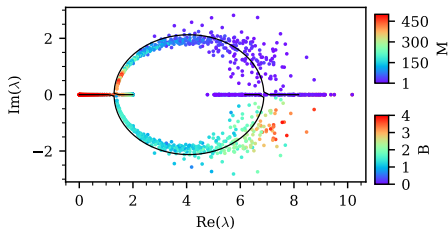
Adding disorder - Gallery

- Spectrum: $\mathbf{W}\psi = -\lambda\psi$

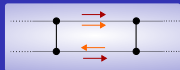
- Four parameters

Propulsion: $\phi_n = \bar{\phi} + [-\sigma_\phi, \sigma_\phi]$

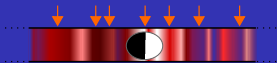
Drift: $f_n = \bar{f} + [-\sigma_f, \sigma_f]$



Participation number: $M = \left[\sum_{n,s} P_{n,s}^2 \right]^{-1}$

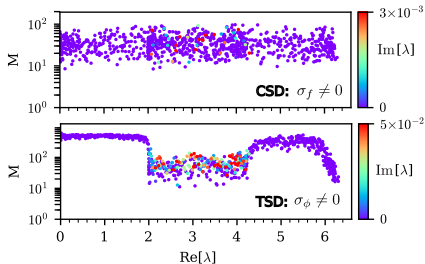
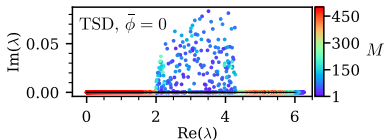


Adding disorder



Adding disorder \Rightarrow Localization ($\sim M^{-1}$)

- CSD (random f_n) \Rightarrow Spectrum is real, localization is uniform (one channel)



- TSD (ϕ_n is random) \Rightarrow Spectrum is complex, no finite threshold for ϕ_n
- Localization drop: One Channel \rightarrow Two channels

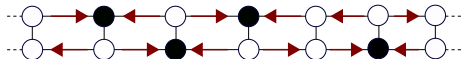
Participation number: $M = \left[\sum_{n,s} P_{n,s}^2 \right]^{-1}$

Floor level

Large TSD: The eigenvalues λ stretch along the real axis.

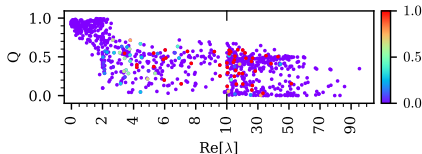
- 25% of the eigenstates stay within the limits $0 < \text{Re}[\lambda] < 2$.
- These eigenvalues remain real.

A floor-level band is formed: symmetric “virtual transitions” occur between the floor sites.



Probability of $|n, s\rangle$ to be in the floor-band:

$$p(\mathcal{E}_{n,s} = \rightarrow) \times p(\mathcal{E}_{n+1,s} = \leftarrow) = \frac{1}{2} \times \frac{1}{2} = \frac{1}{4}$$



Q - Probability to be in the floor-band

Discussion

- Relaxation modes of a stochastic network can be either over-damped or under-damped depending on whether their λ -s are real or complex.
- Without disorder ($\phi_n = \bar{\phi}$), under-damped relaxation require $\bar{\phi} > \phi_c$
- Random ϕ_n - No finite threshold for under-damped relaxation
- Random ϕ_n very different than random f_n (complexity, localization)
- Complexity and de-localization do not come together (contrary to Hatano-Nelson)

[1] DS, Dganit Meidan and Doron Cohen (Phys. Rev. E 98, 012107)