

# The conductance of small mesoscopic disordered rings

Doron Cohen, Ben-Gurion University

## Collaborations:

Tsampikos Kottos (Wesleyan)

Holger Schanz (Gottingen)

Swarnali Bandopadhyay (BGU)

Yoav Etzioni (BGU)

Tal Peer (BGU)

Rangga Budoyo (Wesleyan)

Alex Stotland (BGU)

## Discussions:

Michael Wilkinson (UK)

Bernhard Mehlig (Goteborg)

Yuval Gefen (Weizmann)

Shmuel Fishman (Technion)

\$ISF, \$GIF, \$DIP, \$BSF

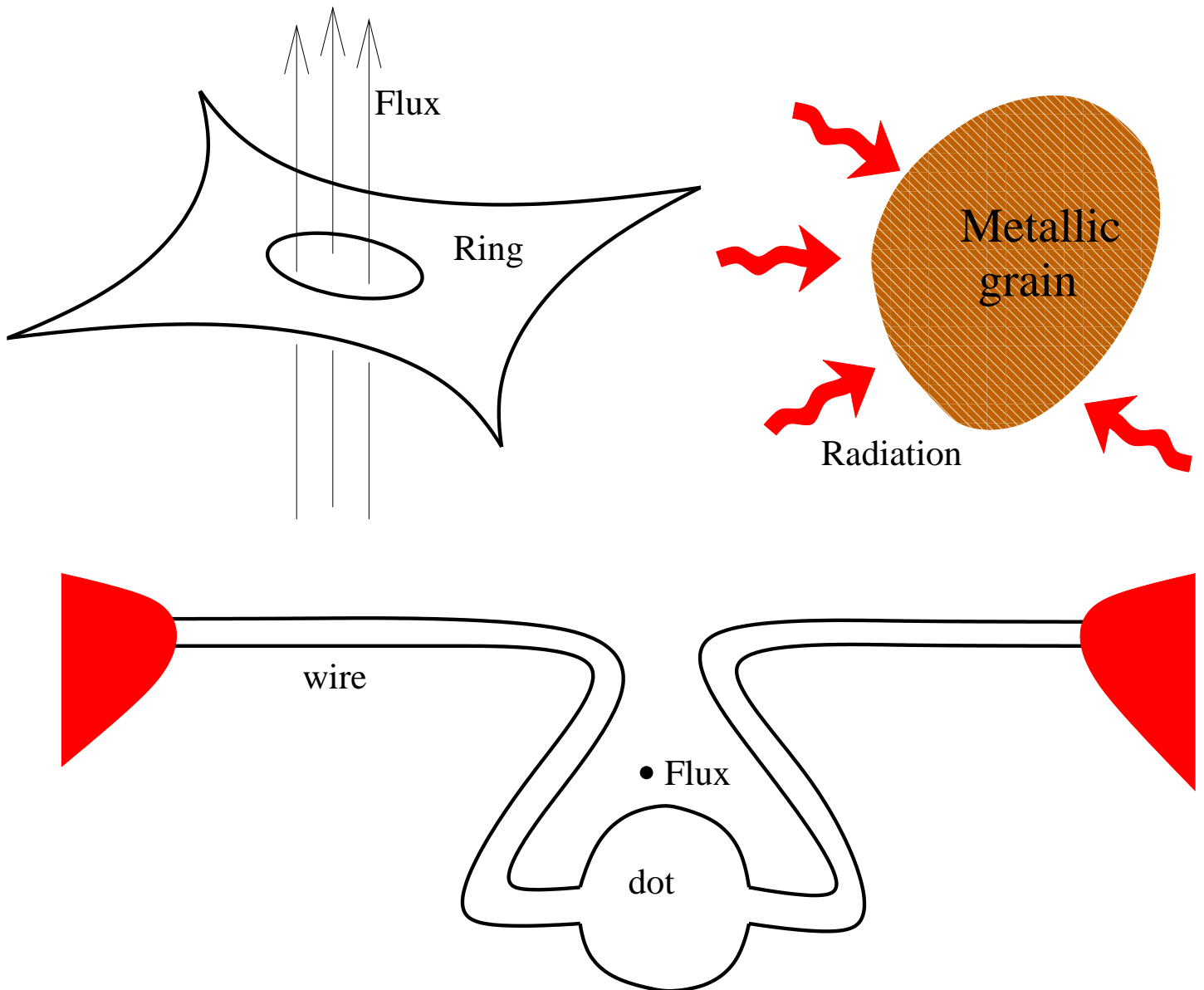
# Driven Systems

Non interacting “spinless” electrons in a ring.

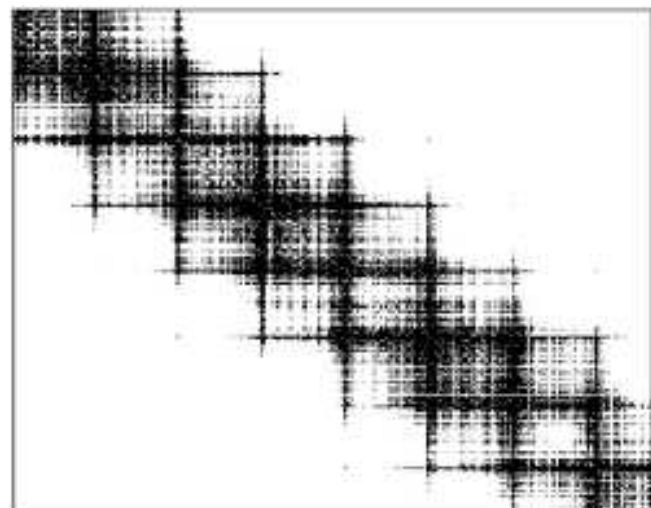
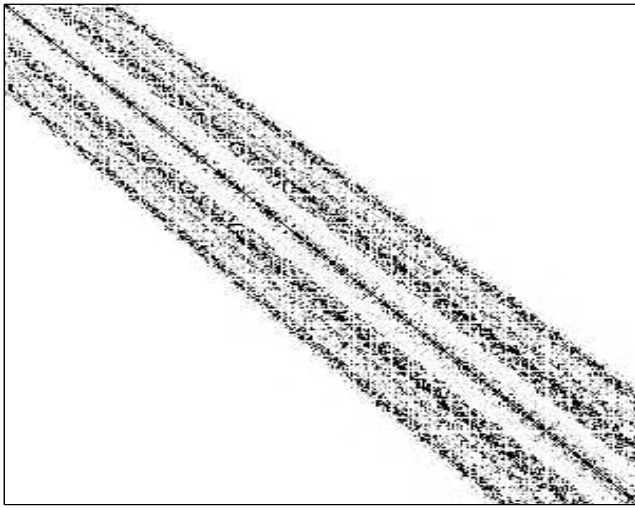
$$\mathcal{H}(Q, P; \Phi(t))$$

$-\dot{\Phi}$  = electro motive force (RMS)

$G \dot{\Phi}^2$  = rate of energy absorption



# Linear Response Theory (LRT)



$$H = \{E_n\} - \Phi(t) \{\mathcal{I}_{nm}\}$$

$$\mathbf{G} = \pi \hbar \sum_{n,m} |\mathcal{I}_{mn}|^2 \delta_T(E_n - E_F) \delta_\Gamma(E_m - E_n)$$

$$\mathbf{G} = \pi \hbar (\varrho(E_F))^2 \langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle$$

applies if

EMF driven transitions  $\ll$  relaxation

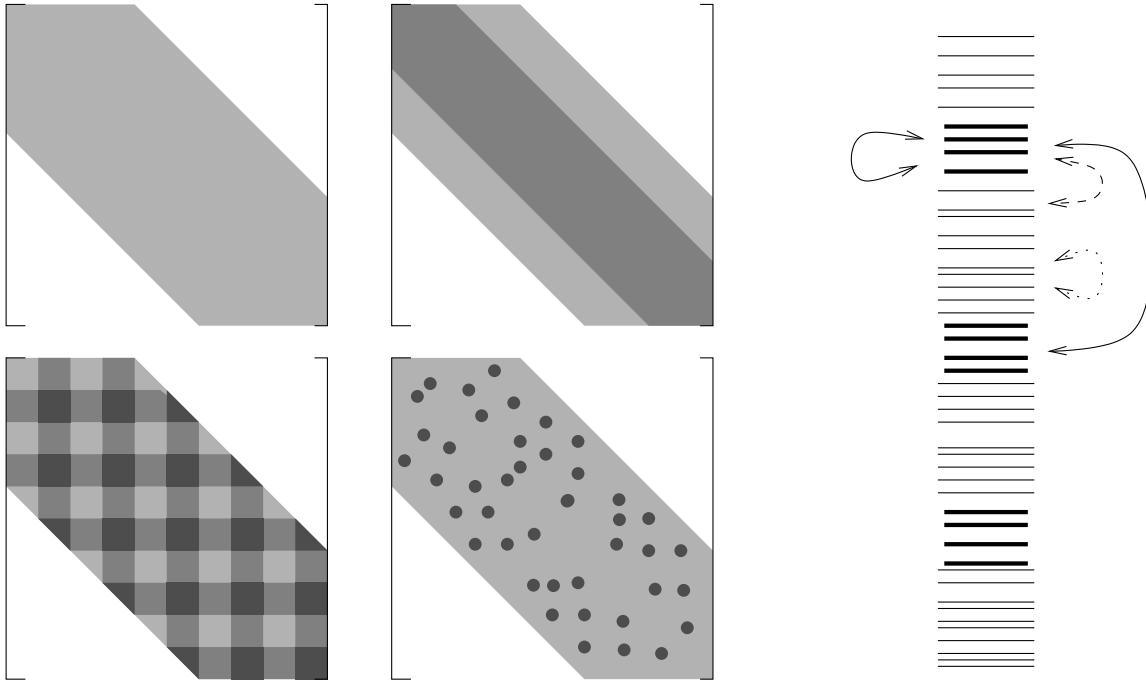
otherwise

*connected sequences of transitions* are essential.

leading to

Semi Linear Response Theory (SLRT)

# Semi Linear Response Theory (SLRT)



$$H = \{E_n\} - \Phi(t) \{I_{nm}\}$$

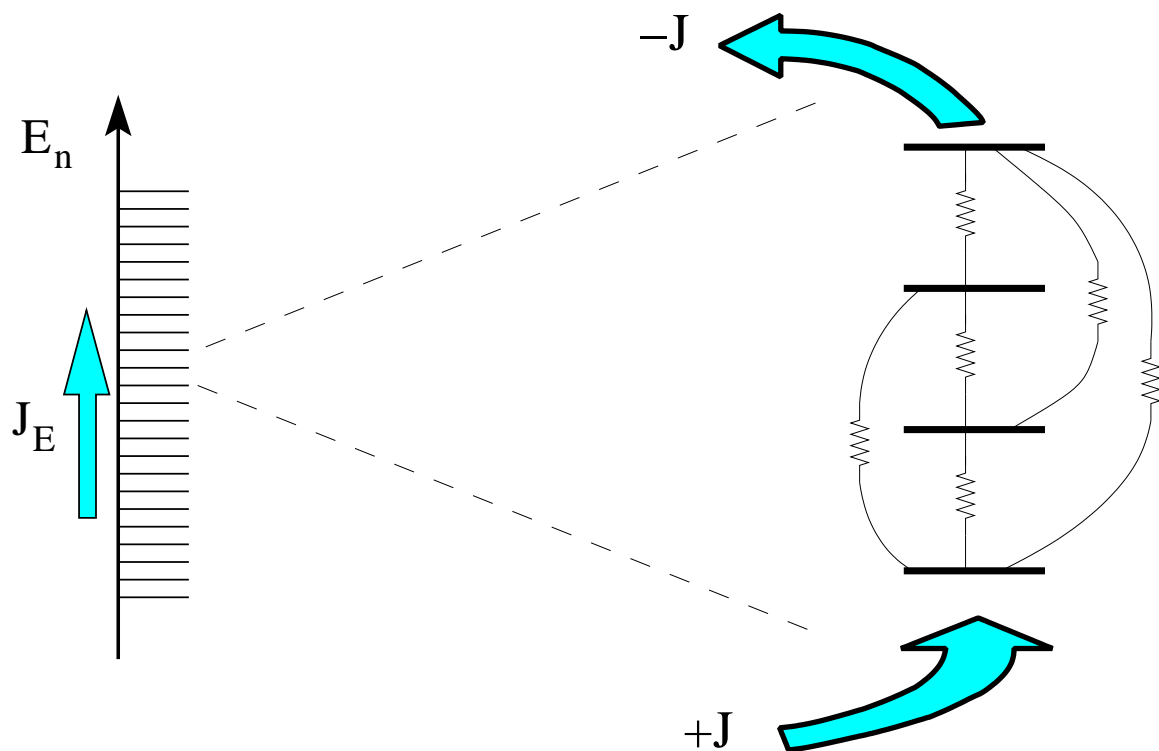
$$\frac{dp_n}{dt} = - \sum_m w_{nm} (p_n - p_m)$$

$$w_{nm} = \text{const} \times g_{nm} \times \text{EMF}^2$$

Scaled transition rates:

$$g_{nm} = 2\rho_F^{-3} \frac{|I_{nm}|^2}{(E_n - E_m)^2} \delta_{\Gamma}(E_n - E_m)$$

## Semi Linear Response Theory (cont.)



$$g_{nm} = 2\rho_F^{-3} \frac{|\mathcal{I}_{nm}|^2}{(E_n - E_m)^2} \delta_{\Gamma}(E_n - E_m)$$

The SLRT analog of the Kubo formula:

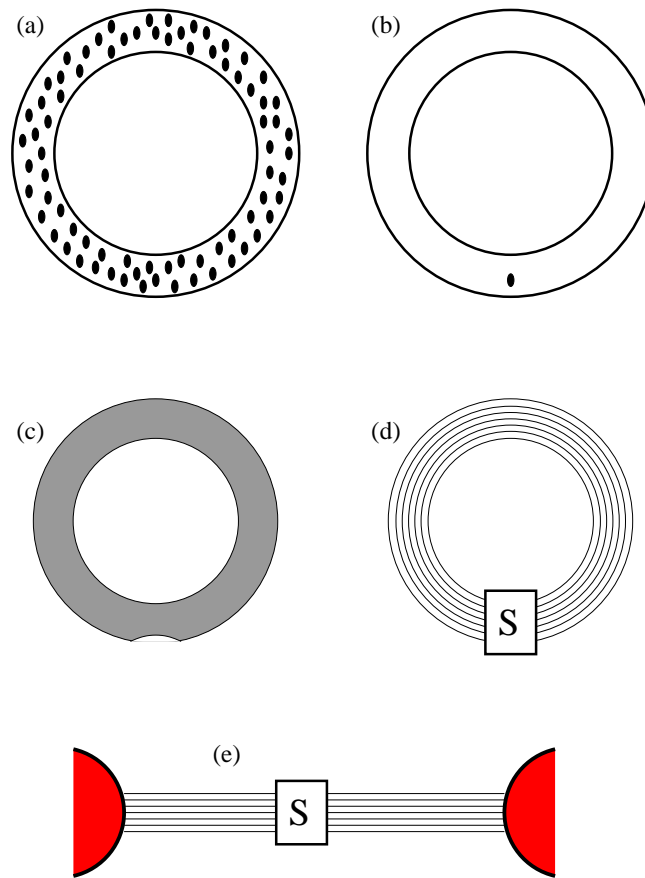
$$\mathbf{G} = \pi\hbar(\rho(E_F))^2 \langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle$$

where

$\langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle \equiv$  inverse resistivity of the network

$$\langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle_{\text{harmonic}} \ll \langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle \ll \langle\langle |\mathcal{I}_{mn}|^2 \rangle\rangle_{\text{algebraic}}$$

# Conductance of mesoscopic rings



Naive expectation (assuming  $\Gamma > \Delta$ ):

$$G = \frac{e^2}{2\pi\hbar} \mathcal{M} \frac{\ell}{L} + \mathcal{O}\left(\frac{\Delta}{\Gamma}\right)$$

$L$  = perimeter of the ring

$\ell$  = mean free path  $\propto W^2$

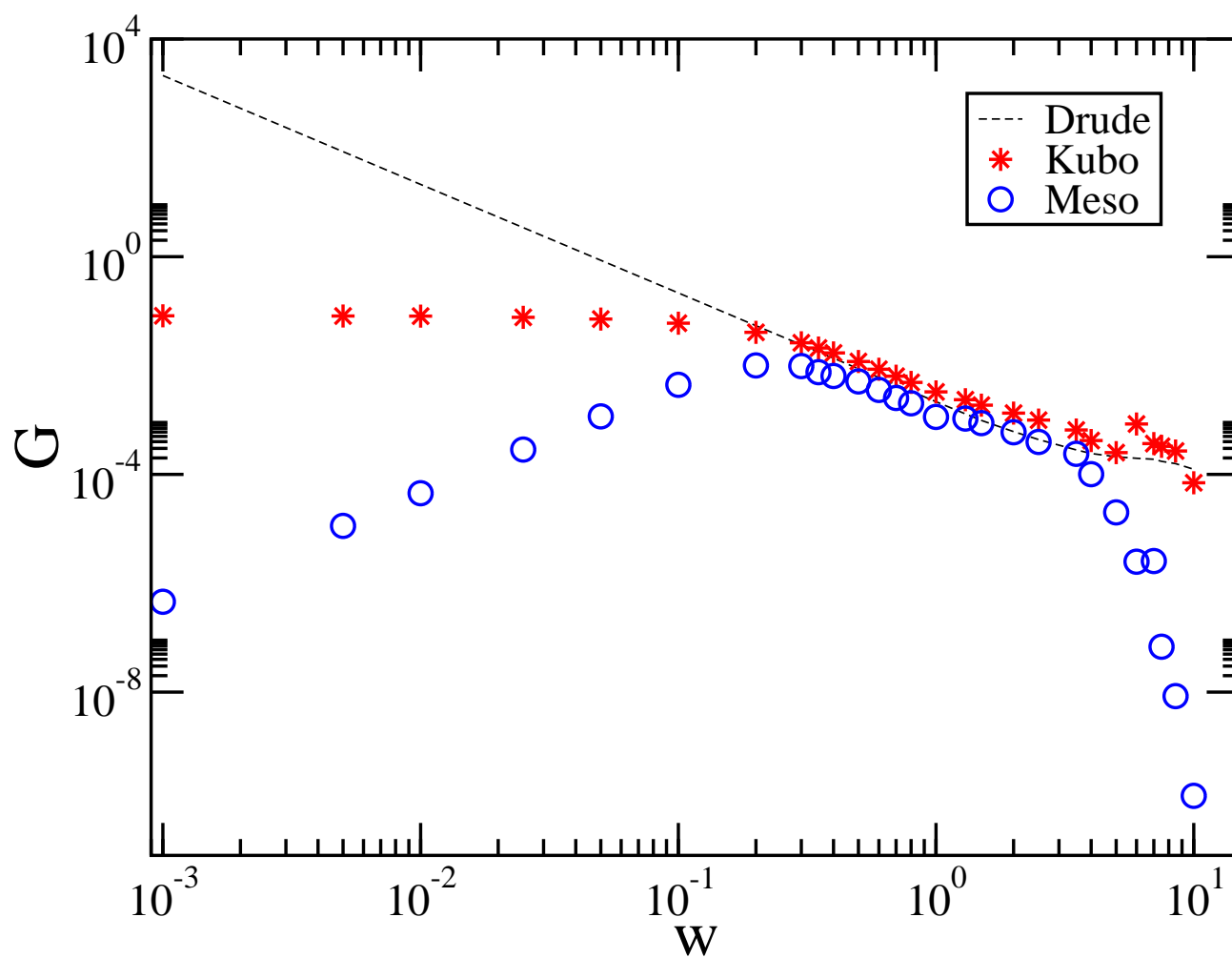
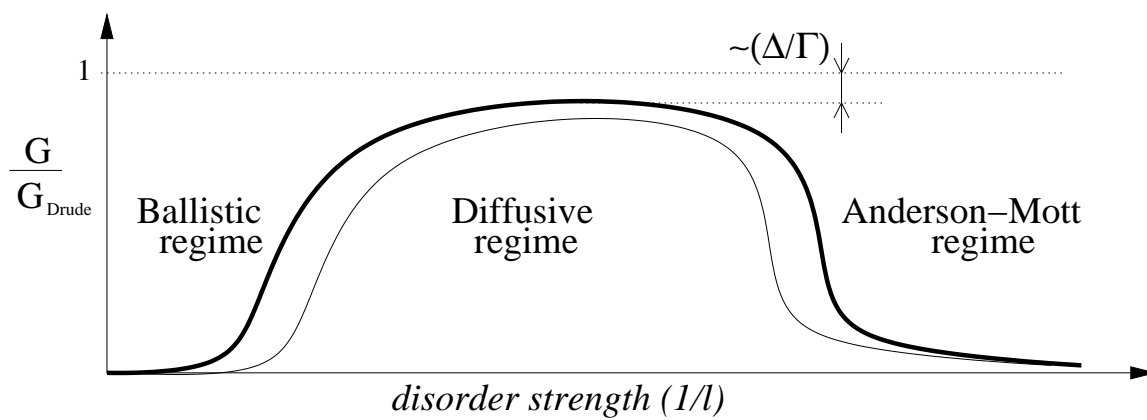
$l_\infty$  = localization length  $\approx \mathcal{M}\ell$

Ballistic regime:  $L \ll \ell$

Diffusive regime:  $\ell \ll L \ll l_\infty$

Anderson regime:  $l_\infty \ll L$

# Conductance versus disorder

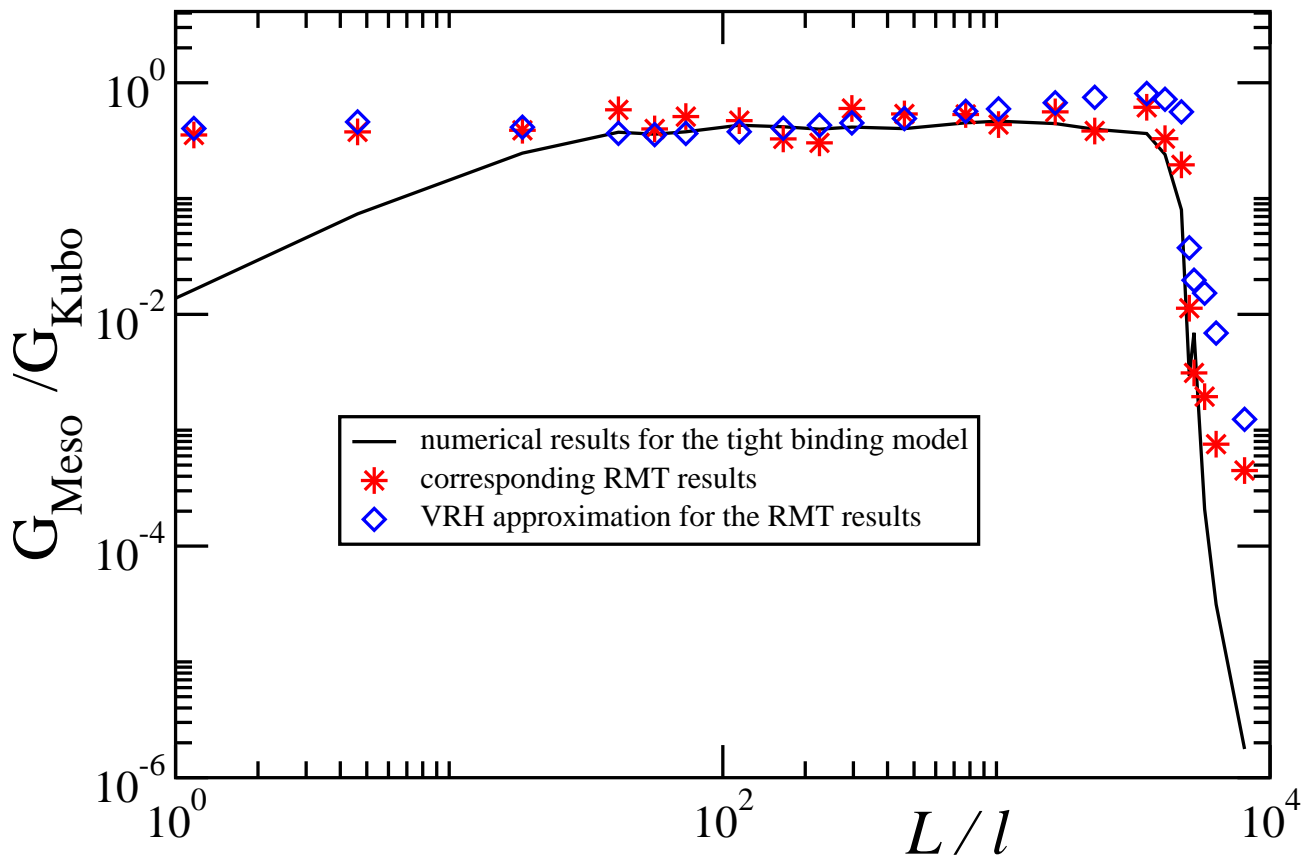


# The RMT modeling

$$\{|v_{nm}|^2\} \equiv \{X\}$$

Characterization of the perturbation matrix:

- bandwidth ( $b$ )
- sparsity ( $p$ )
- texture



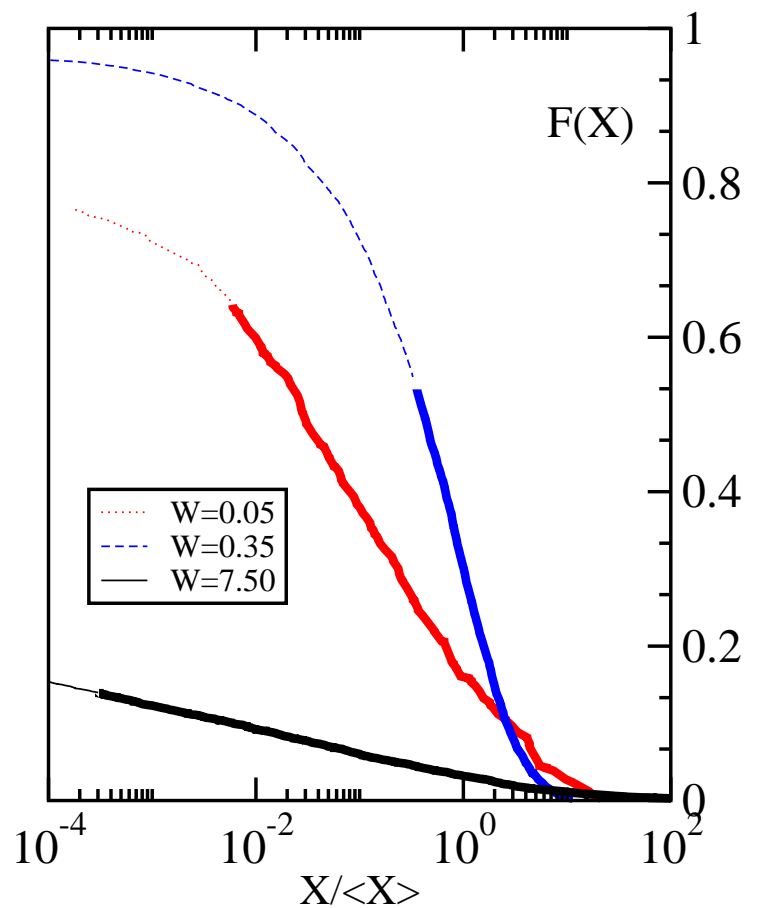
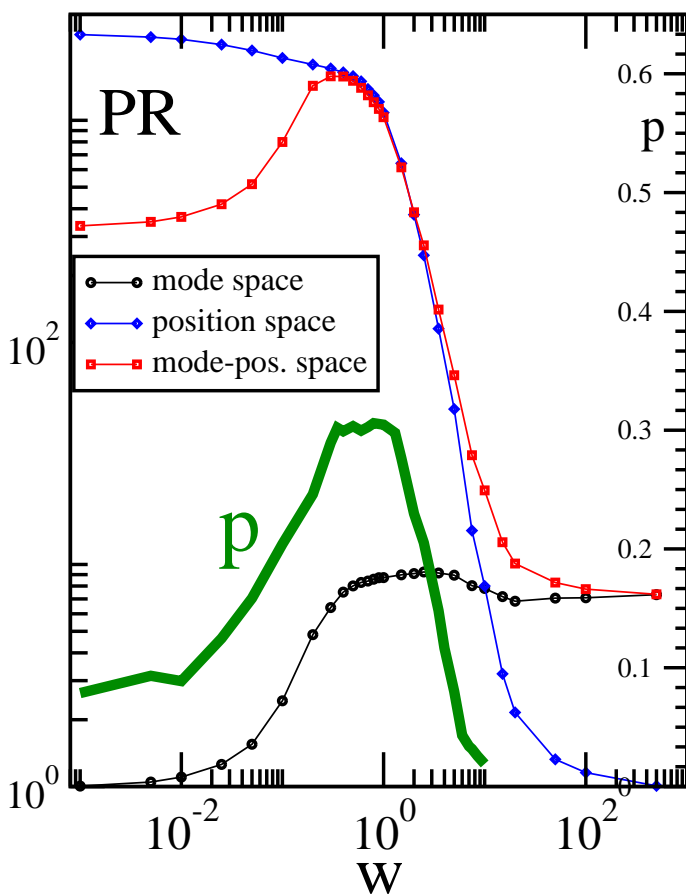
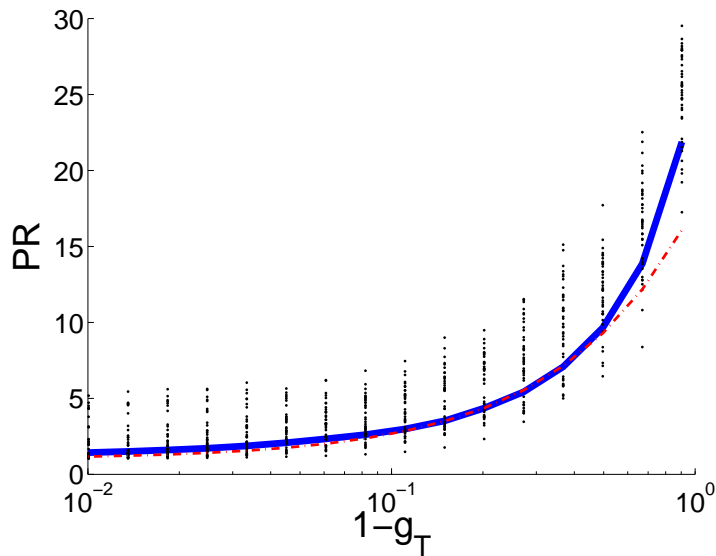
Comparison between:

- Actual results based on “real” matrices
- RMT results based on “artificial” matrices
- Semi-analytical VRH estimate



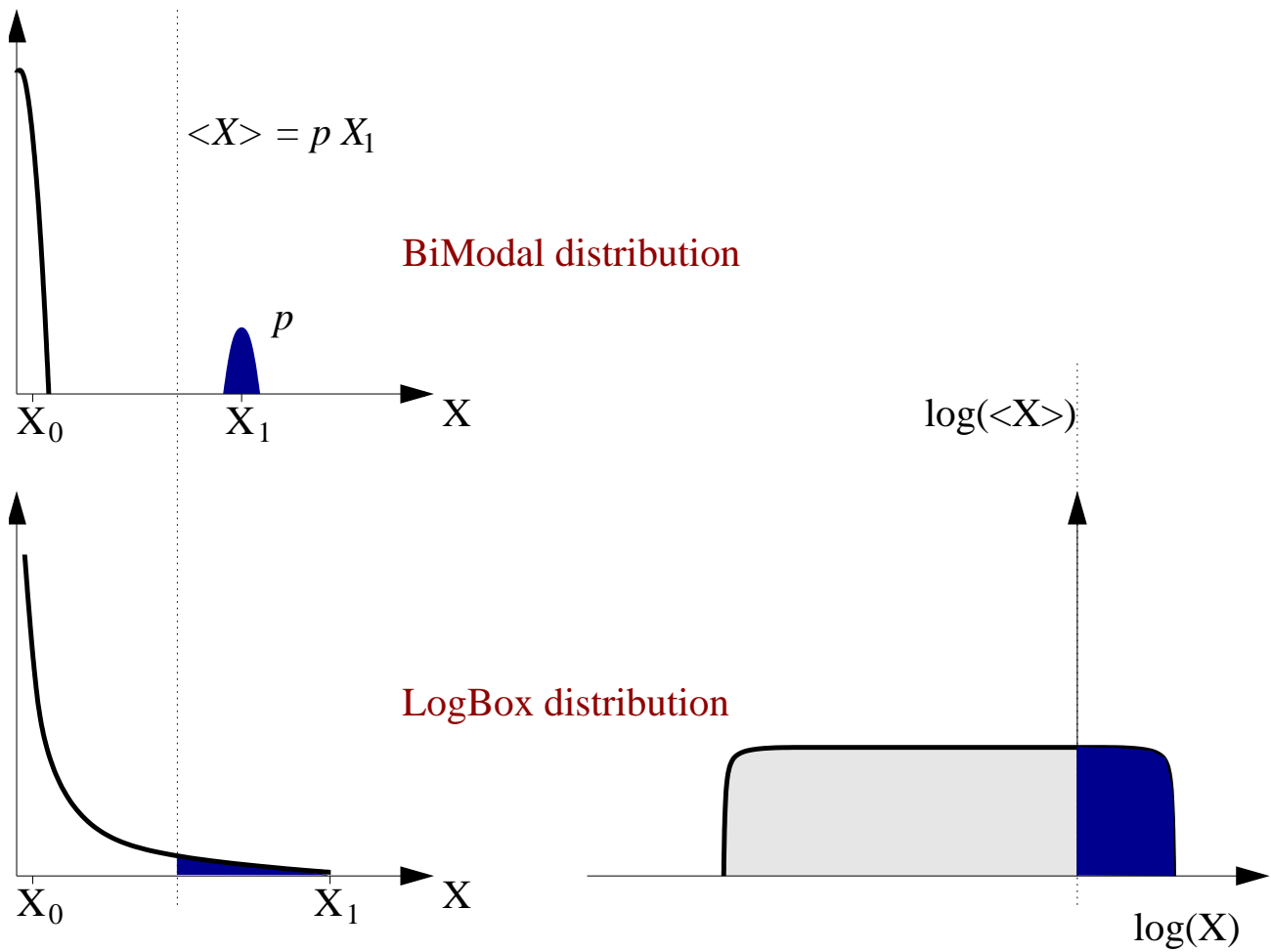
# Ergodicity of the eigenstates

- **Weak disorder** (ballistic rings):  
Wavefunctions are localized in mode space.
- **Strong disorder** (Anderson localization):  
Wavefunctions are localized in real space.



# Modeling of sparsity

$$X \equiv |v_{nm}|^2 \sim \frac{1}{\mathcal{M}^2} v_F^2 \exp\left(-\frac{x}{l_\infty}\right)$$



$$X \in \text{LogBox}[X_0, X_1]$$

$$\tilde{p} \equiv (\ln(X_1/X_0))^{-1}$$

$$p \equiv \text{Prob}(X > \langle X \rangle) \approx -\tilde{p} \ln \tilde{p}$$

$$\langle X \rangle \approx \tilde{p} X_1 \sim p X_1$$

## The VRH estimate

$$\mathbf{G} = \pi \hbar \left( \frac{e}{L} \right)^2 \sum_{n,m} |v_{mn}|^2 \delta_T(E_n - E_F) \delta_{\Gamma}(E_m - E_n)$$

$$\mathbf{G} = \frac{1}{2} \left( \frac{e}{L} \right)^2 \rho_F \int \tilde{C}_{\text{qm}}(\omega) \delta_{\Gamma}(\omega) d\omega$$

$$\tilde{C}_{\text{qm-LRT}}(\omega) \equiv 2\pi \rho_F \langle X \rangle$$

$$\tilde{C}_{\text{qm-SLRT}}(\omega) \equiv 2\pi \rho_F \bar{X}$$

where by definition:  $\left( \frac{\omega}{\Delta} \right) \text{Prob}(X > \bar{X}) \sim 1$

For strong disorder we get:

$$\bar{X} \approx v_F^2 \exp\left(-\frac{\Delta \ell}{\omega}\right)$$

$$\mathbf{G} \propto \int \exp\left(-\frac{\Delta \ell}{|\omega|}\right) \exp\left(-\frac{|\omega|}{\omega_c}\right) d\omega$$

# LRT, SLRT and beyond

$-\dot{\Phi}$  = electro motive force (RMS)

$G \dot{\Phi}^2$  = rate of energy absorption

## Semi linear response theory

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*Beyond FGR*,  
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