

BEC dynamics in few site systems

Maya Chuchem [1,3]

Ben-Gurion University

Collaborations:

Doron Cohen (BGU/Phys) [1,2,3]

Tsampikos Kottos (Wesleyan) [3]

Katrina Smith-Mannschott

(Wesleyan / Gottingen) [3]

Moritz Hiller (Freiburg) [3]

Amichay Vardi (BGU/Chem) [1,2,3]

Erez Boukobza (BGU/Chem) [1,2]

[1] Occupation dynamics & fluctuations (PRL 2009)

[2] Dynamical phase locking of BECs (PRA 2009)

[3] Semiclassical analysis & quantum dynamics (arXiv 2010)

Thanks:

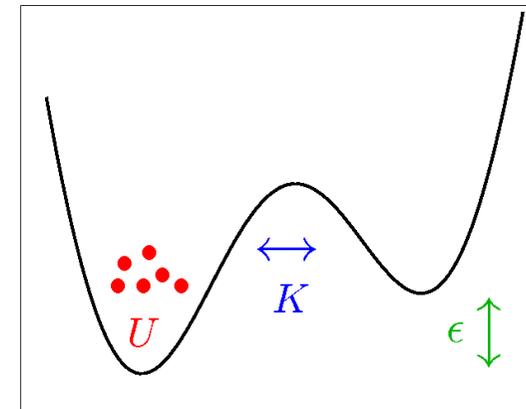
Isaac Israel

Alexander Stotland

Itamar Sela

Yoav Etzioni

\$BSF, \$DIP, \$FOR760



The model

The Bose-Hubbard Hamiltonian (BHH) for a dimer

$$\mathcal{H} = \sum_{i=1,2} \left[\mathcal{E}_i \hat{n}_i + \frac{U}{2} \hat{n}_i (\hat{n}_i - 1) \right] - \frac{K}{2} (\hat{a}_2^\dagger \hat{a}_1 + \hat{a}_1^\dagger \hat{a}_2)$$

K = hopping

U = interaction

$\mathcal{E} = \mathcal{E}_2 - \mathcal{E}_1 =$ bias

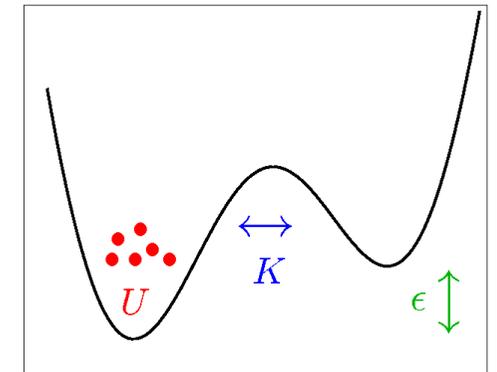
N particles in a double well is like spin $j = N/2$ system

$$\mathcal{H} = -\mathcal{E} \hat{J}_z + U \hat{J}_z^2 - K \hat{J}_x + \text{const}$$

$$\hat{J}_z = (N/2) \cos(\theta) = \underline{\hat{n} = \text{occupation difference}}$$

$$\hat{J}_x \approx (N/2) \sin(\theta) \cos(\varphi), \quad \underline{\varphi = \text{relative phase}}$$

$$(J_x, J_y, J_z) \Leftrightarrow (\theta, \varphi) \Leftrightarrow (n, \varphi)$$



$$u \equiv \frac{NU}{K}, \quad \varepsilon \equiv \frac{\mathcal{E}}{K}$$

Phase space analysis

$$\mathcal{H} = -\varepsilon \hat{J}_z + U \hat{J}_z^2 - K \hat{J}_x \quad u \equiv \frac{NU}{K}, \quad \varepsilon \equiv \frac{\mathcal{E}}{K}$$

$$\mathcal{H}(\theta, \varphi) = \frac{NK}{2} \left[\frac{1}{2} u (\cos \theta)^2 - \varepsilon \cos \theta - \sin \theta \cos \varphi \right] \quad \text{spherical phase space}$$

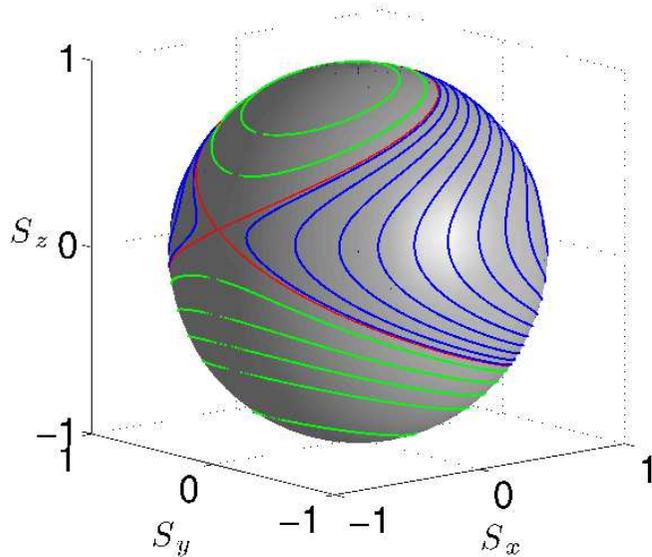
$$\mathcal{H}(\hat{n}, \varphi) = (\text{similar to Josephson/pendulum Hamiltonian}) \quad \text{cylindrical}$$

For $u > 1$ and $|\varepsilon| < \varepsilon_c$

Sea, Islands, Separatrix

$$\varepsilon_c = (u^{2/3} - 1)^{3/2}$$

$$A_c \approx 4\pi (1 - u^{-2/3})^{3/2}$$



Rabi regime: $u < 1$ (no islands)

Josephson regime: $1 < u < N^2$

Fock regime: $u > N^2$ (empty sea)

WKB quantization

$$h = \text{Planck cell area in steradians} = \frac{4\pi}{N+1}$$

$$A(E_n) = \left(\frac{1}{2} + n\right) h$$

$$\omega(E) \equiv \frac{dE}{dn} = \left[\frac{1}{h} A'(E)\right]^{-1}$$

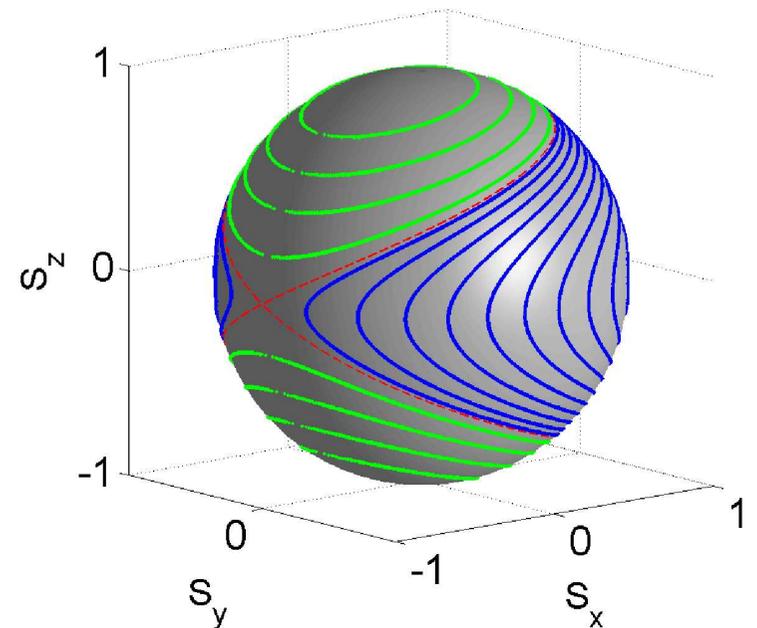
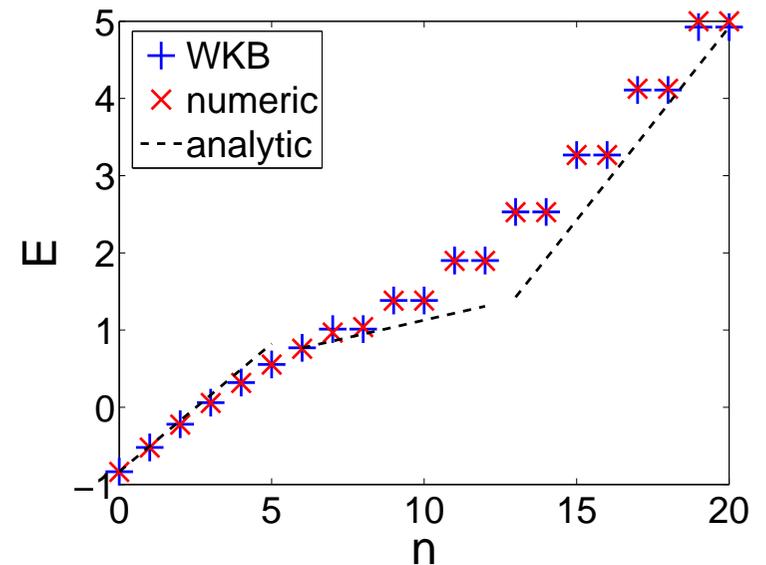
$$\omega_K \approx K = \text{Rabi Frequency}$$

$$\omega_J \approx \sqrt{NUK} = \sqrt{u} \omega_K$$

$$\omega_+ \approx NU = u \omega_K$$

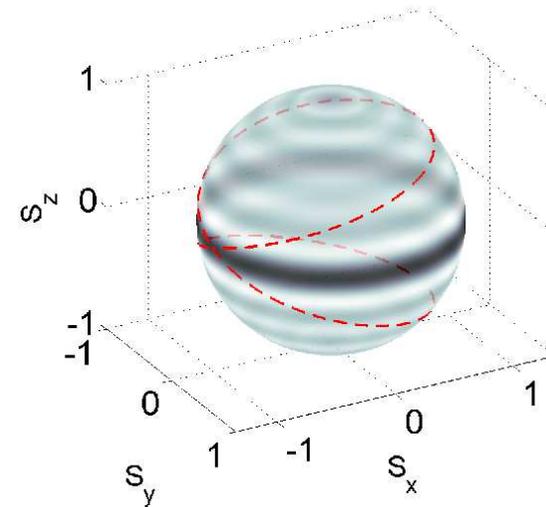
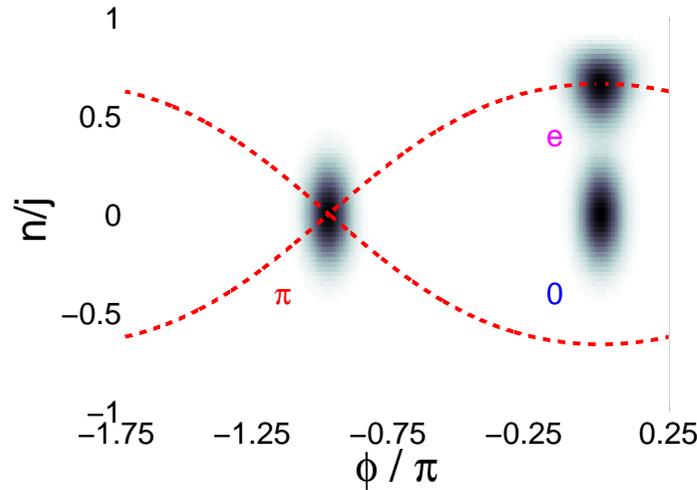
$$\omega_x \approx \left[\frac{1}{2} \log\left(\frac{N^2}{u}\right)\right]^{-1} \omega_J$$

$$E_{-,x} = \mp(1/2)NK, \quad E_+ = (1/4)[u+(1/u)]NK$$



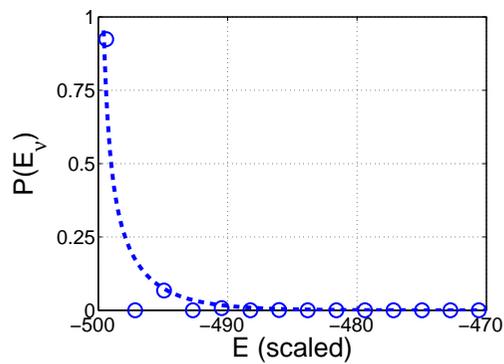
Typical preparations and their LDOS

Coherent state
 $|\theta\varphi\rangle$
 is like a
 minimal
 Gaussian
 wavepacket



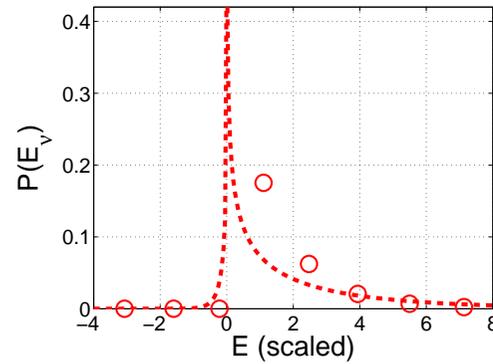
Fock state
 $|n\rangle$
 is like
 equi-latitude
 annulus

Zero preparation '0'



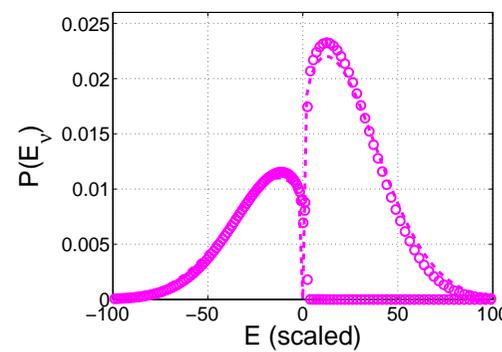
$$P(E) \sim \text{BesselI} \left[\frac{E - E_-}{NU} \right]$$

Pi preparation 'pi'



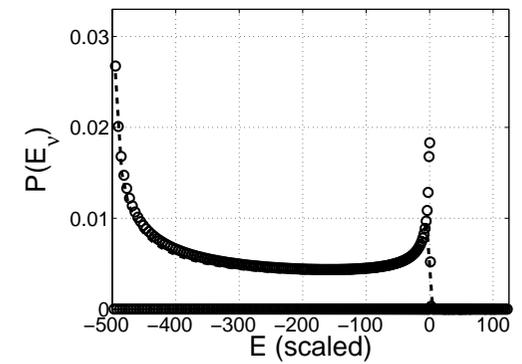
$$\sim \frac{\omega_x}{\omega_J} \text{BesselK} \left[\frac{E - E_x}{NU} \right]$$

Edge preparation 'e'



$$\sim \frac{\omega(E)}{\omega_J} \exp \left[-\frac{1}{N} \left(\frac{E - E_x}{\omega_J} \right)^2 \right]$$

TwinFock preparation



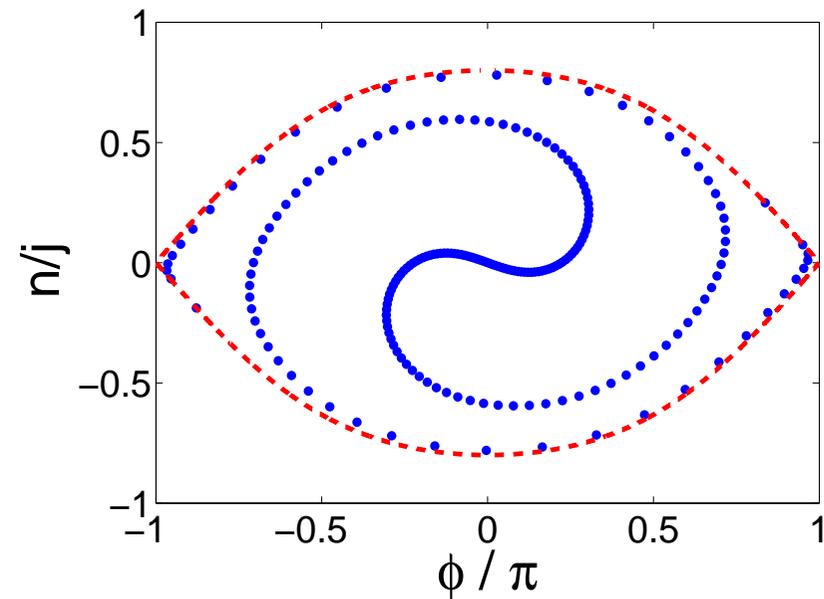
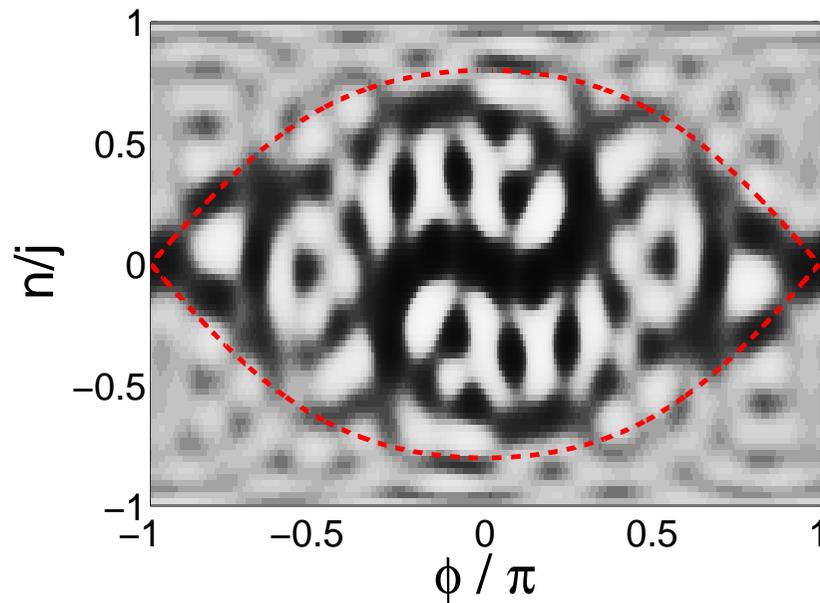
$$\sim \frac{\omega(E)}{\omega_J} \left[1 - \left(\frac{2E}{NK} \right)^2 \right]^{-1/2}$$

Numerical simulations - analysis of the evolution

MeanField theory (GPE) = classical evolution of a **point** in phase space

SemiClassical theory = classical evolution of a **distribution** in phase space

Quantum theory = recurrences, fluctuations (WKB is very good)



Any operator \hat{A} can be presented by the phase-space function $A_W(\Omega) = \text{trace}[\hat{P}^\Omega \hat{\rho}]$

where the projector-like operators obey: $\int \frac{d\Omega}{h} \hat{P}^\Omega = \hat{\mathbf{1}}$ and for Hermitian operators $A_W \in \mathcal{R}$

The expectation value: $\langle \hat{A} \rangle = \text{trace}[\hat{\rho} \hat{A}] = \int \frac{d\Omega}{h} \rho_W(\Omega) A_W(\Omega)$

Recurrences and fluctuations

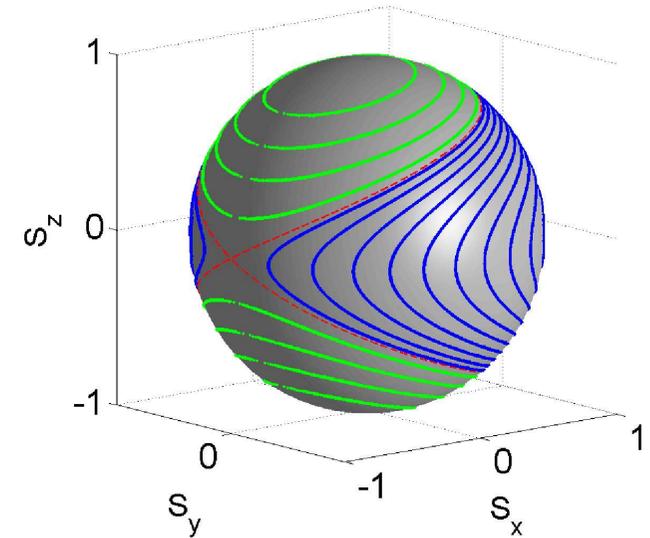
$$\vec{S} = \langle \vec{J} \rangle / j \quad , \quad j = \frac{N}{2}$$

$$\text{OccupationDiff} = (N/2) [1 + S_z]$$

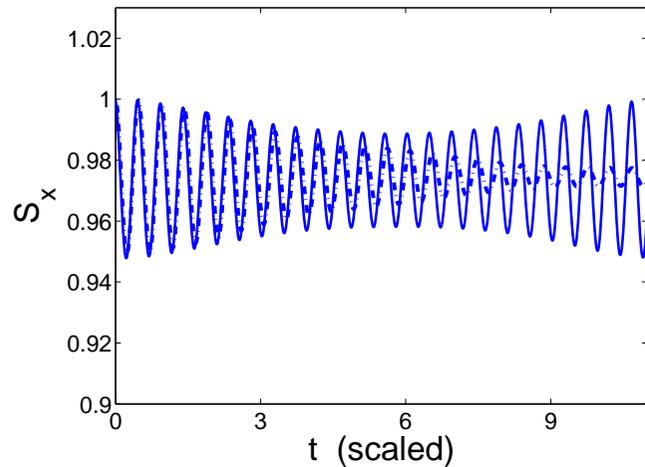
$$\text{OneBodyPurity} = (1/2) [1 + S_x^2 + S_y^2 + S_z^2]$$

$$\text{FringeVisibility} = [S_x^2 + S_y^2]^{1/2}$$

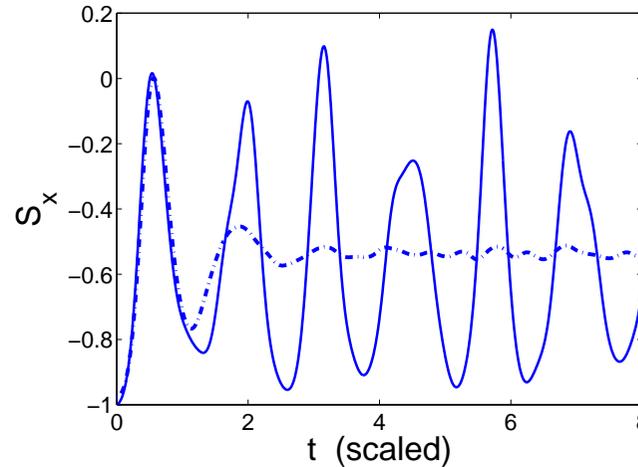
$$\text{RelativePhase} = \arctan(S_x/S_y)$$



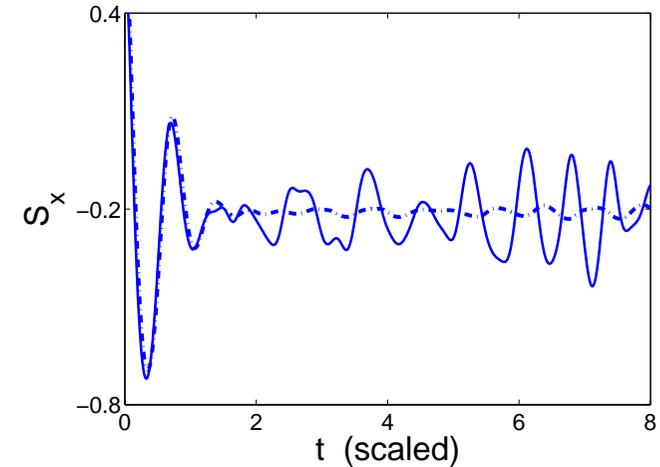
Zero prep



Pi prep

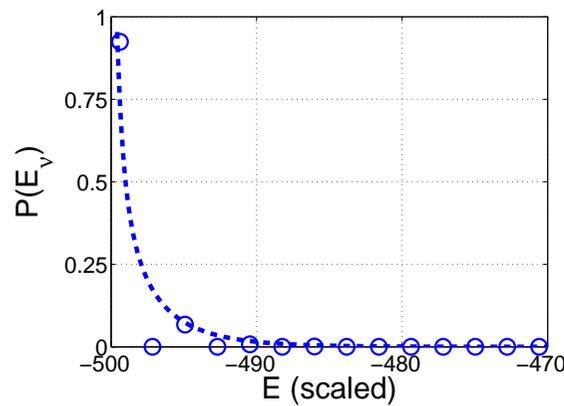
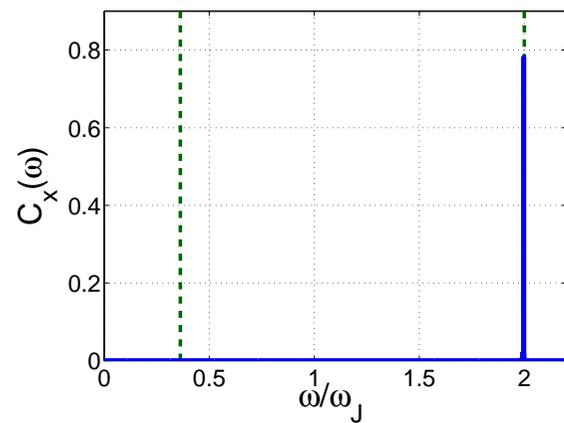
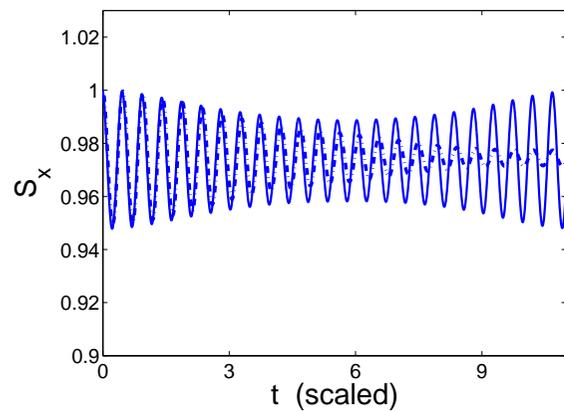


Edge prep

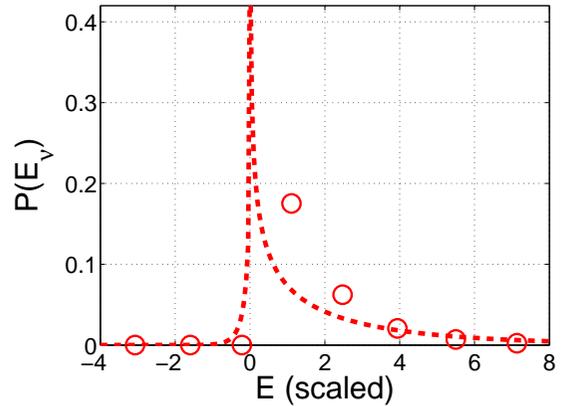
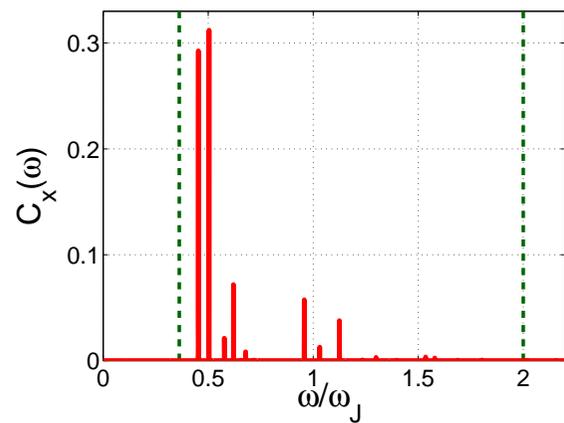
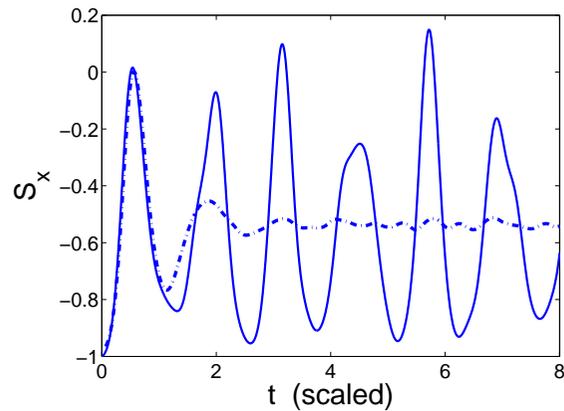


Temporal behavior & the spectral content of S_x

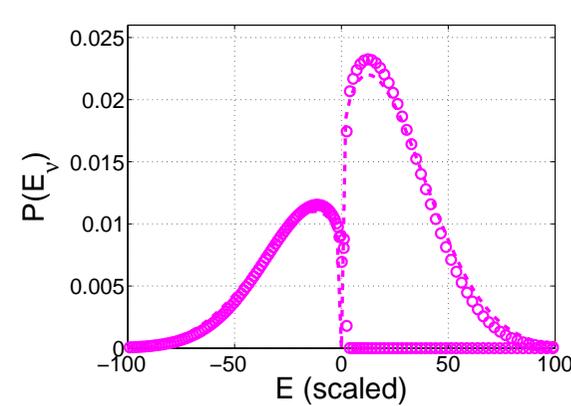
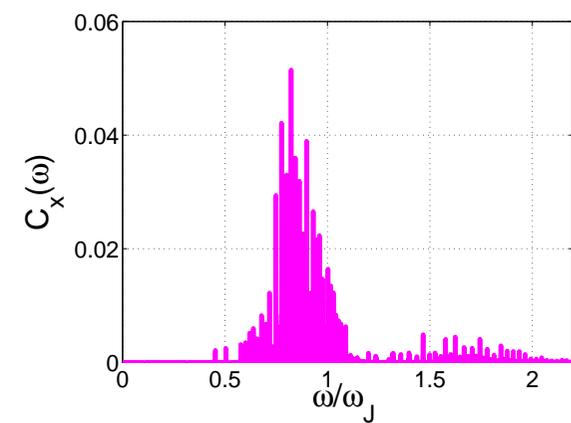
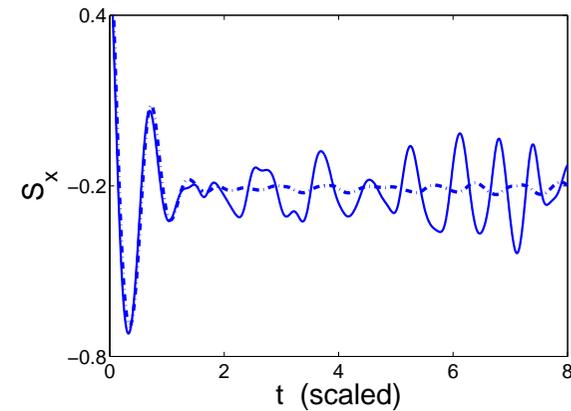
Zero



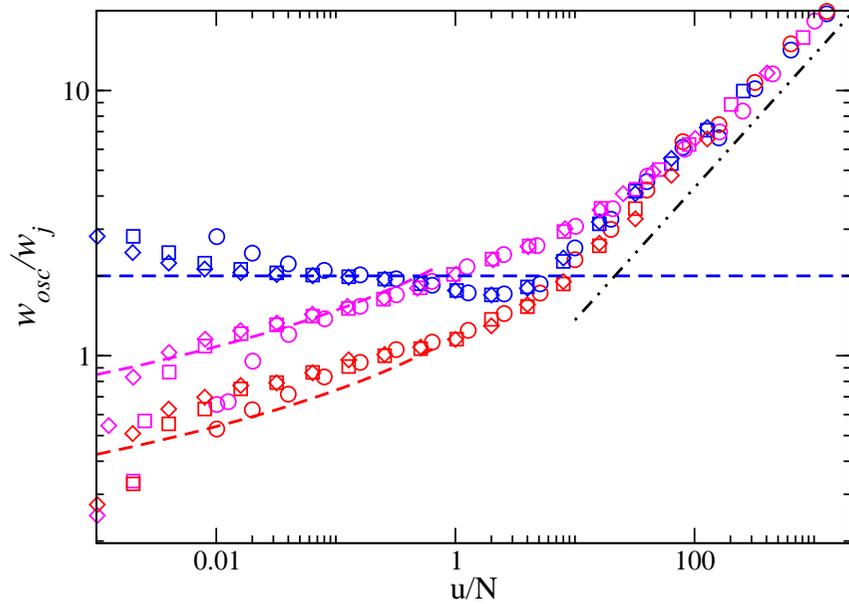
Pi



Edge



The spectral content of S_x

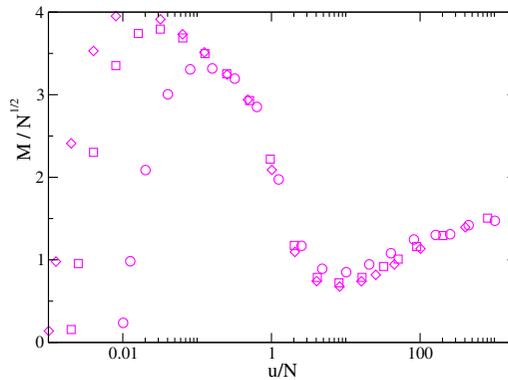
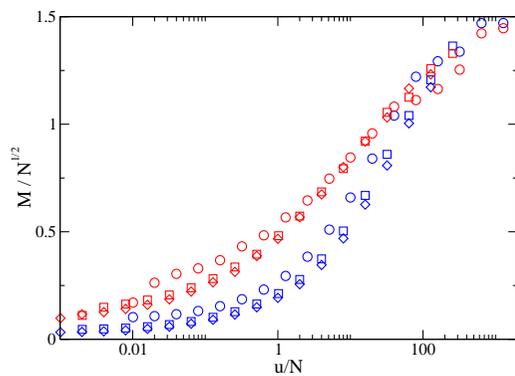


$$\omega_{\text{osc}} \approx 2 \left[\log \left(\frac{N}{u} \right) \right]^{-1} \omega_J \quad [\text{Edge}]$$

$$\omega_{\text{osc}} \approx \left[\log \left(\frac{N}{u} \right) \right]^{-1} \omega_J \quad [\text{Pi}]$$

$$\omega_{\text{osc}} \approx \left(\frac{u}{N} \right)^{1/2} \omega_J \quad [u \gg N]$$

The semiclassical ratio $\sim (u/N)^{1/2}$



The participation number:

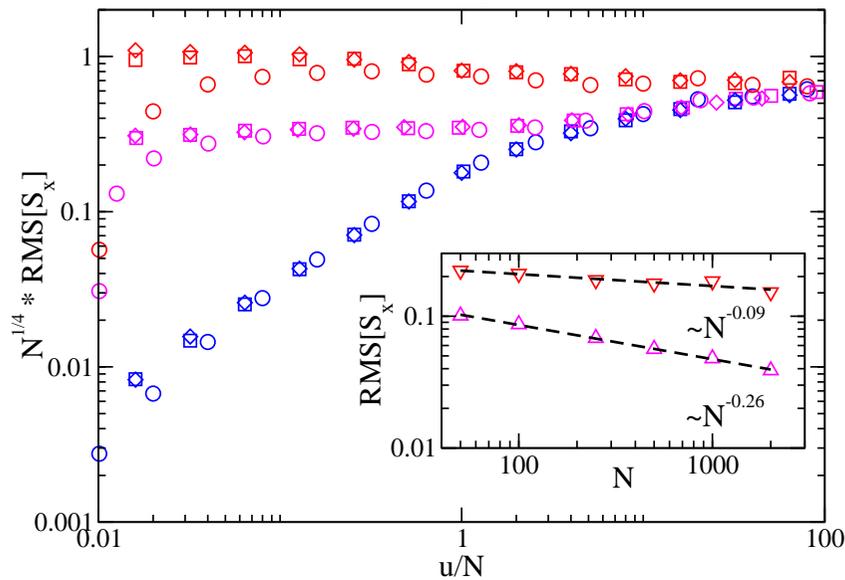
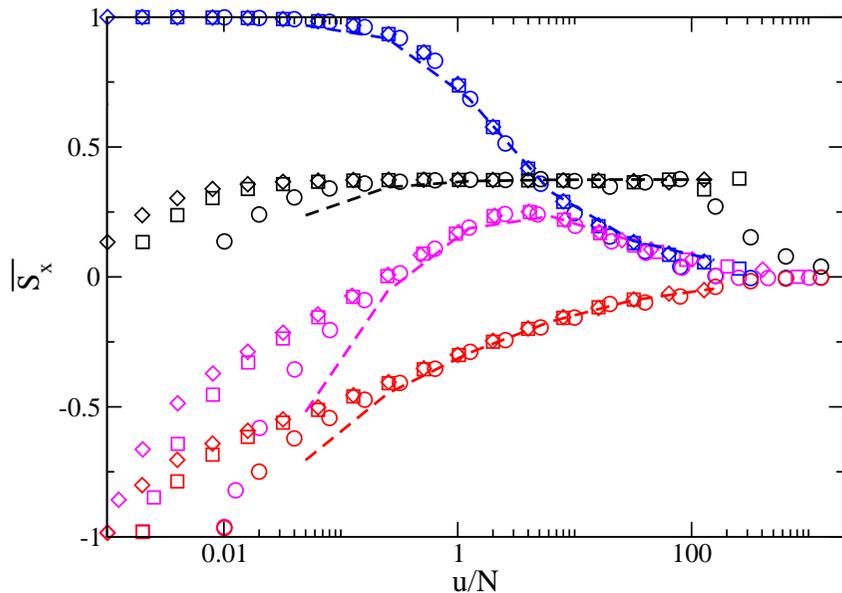
$$M = \left[\sum_{\nu} P(E_{\nu})^2 \right]^{-1}$$

$$M \approx \left[\log \left(\frac{N}{u} \right) \right] \sqrt{N} \quad [\text{Edge}]$$

$$M \approx \left[\log \left(\frac{N}{u} \right) \right] \sqrt{u} \quad [\text{Pi}]$$

“Zero” , “Pi” , “Edge” . $N = 100$ \circ , $N = 500$ \square , $N = 1000$ \diamond .

Fluctuations of S_x



“Zero”, “Pi”, “Edge”, “TwinFock” .

$N = 100$ \circ , $N = 500$ \square , $N = 1000$ \diamond .

$$\overline{S_x} \approx \exp\left[-\frac{1}{4}(u/N)\right] \quad [\text{Zero}]$$

$$\overline{S_x} \approx -1 - 4/\log\left[\frac{1}{32}(u/N)\right] \quad [\text{Pi}]$$

$$\overline{S_x} \approx 1/3 \quad [\text{TwinFock}]$$

$$\text{RMS} [\langle A \rangle_t] = \left[\frac{1}{M} \int \tilde{C}_{\text{cl}}(\omega) d\omega \right]^{1/2}, \quad \left[\frac{M}{\sqrt{N}} \sim f\left(\sqrt{\frac{u}{N}}\right) \right]$$

$$\text{RMS} [S_x(t)] \sim N^{-1/4} \quad [\text{Edge}]$$

$$\text{RMS} [S_x(t)] \sim (\log(N))^{-1/2} \quad [\text{Pi}]$$

“Zero” Coherence maintained if $u/N < 1$ (phase locking).

“Pi” Fluctuations are suppressed by u .

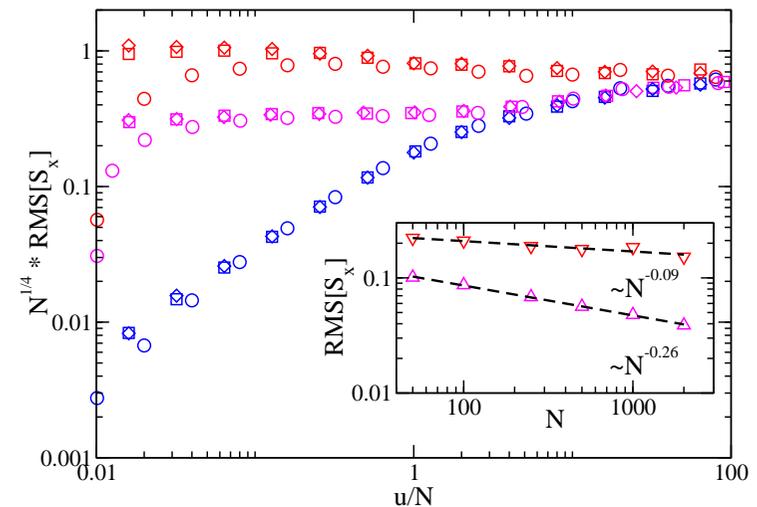
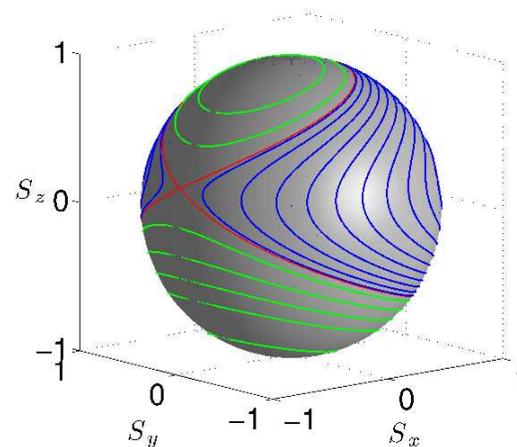
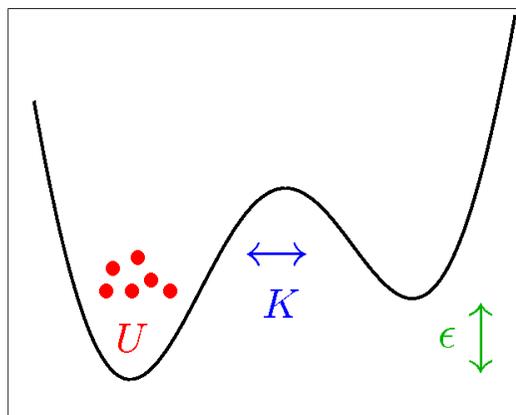
“Edge” Fluctuations are suppressed by N (classical limit).

“TwinFock” Self induced coherence leading to $\overline{S_x} \approx 1/3$.

The semiclassical ratio $\sim (u/N)^{1/2}$

Summary

- One-particle coherence loss and buildup analysis using semiclassical phase-space picture.
- The simplicity of BHH allows semi-analytic WKB quantization.
- Closed semiclassical results for LDOS provide useful insights for the quantum evolution.
- The long-time fringe visibility of an initially coherent state in the Josephson interaction regime, has a u/N (semi-classical ratio) dependent value.
- The Zero prep remains roughly Gaussian throughout its motion, a Pi prep squeezes rapidly and its relative-phase information is lost.
- Two types of coherent preparations in the vicinity of the separatrix show significant differences in their M dependence on $(u; N)$.
- The Pi prep preparation exhibits u dependent fluctuations, whereas the Edge prep exhibits N dependent fluctuations.



Thank You